

# Amplitude and Importance of Relativistic Contributions to the Galaxy Bispectrum

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**Relativistic Effects and Novel Observables in Cosmology**

**University of Geneva**

**July 2024**



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# Decoupling Local Primordial non-Gaussianity from Relativistic Effects in the Galaxy Bispectrum

Rossiter *et al.* (2024) - tomorrow!

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# Expert Guidance



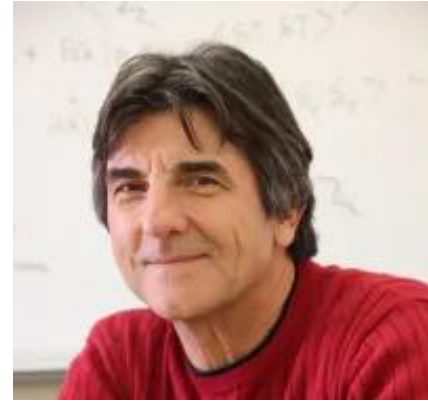
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**Stefano Camera**  
University of Turin



**Chris Clarkson**  
Queen Mary University  
of London



**Roy Maartens**  
University of the  
Western Cape

# Background papers



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Local primordial non-Gaussianity in  
the relativistic galaxy bispectrum

[arXiv:2011.13660](https://arxiv.org/abs/2011.13660)

Roy Maartens<sup>1,2</sup>, Sheean Jolicoeur<sup>1</sup>, Obinna Umeh<sup>2</sup>,  
Eline M. De Weerd<sup>3</sup>, Chris Clarkson<sup>3,1</sup>

Detecting the relativistic galaxy  
bispectrum

[arXiv:1911.02398](https://arxiv.org/abs/1911.02398)

Roy Maartens<sup>1,2</sup>, Sheean Jolicoeur<sup>1</sup>, Obinna Umeh<sup>2</sup>,  
Eline M. De Weerd<sup>3</sup>, Chris Clarkson<sup>3,1,4</sup>, Stefano Camera<sup>5,6,1</sup>

# Contributions to the Galaxy Bispectrum



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- The Universe is not Gaussian: Non-linear evolution of structures and gravitational dynamics, Primordial non-Gaussianity, etc.
- Require higher order statistics such as the Bispectrum

$$\langle \delta_g(\mathbf{k}_1) \delta_g(\mathbf{k}_2) \delta_g(\mathbf{k}_3) \rangle = (2\pi)^3 B_g(\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3) \delta^D(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3)$$

## Relativistic effects and local primordial non-Gaussianity

- Source of non-Gaussianity in the observed galaxy distribution
- More accessible in the bispectrum than in the power spectrum - long modes and short modes couple
- Imaginary terms don't become real when using a single tracer - Smoking gun dipole

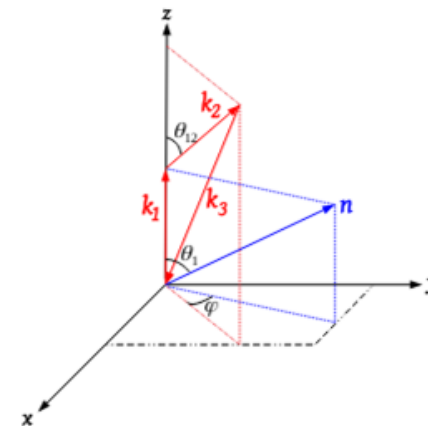
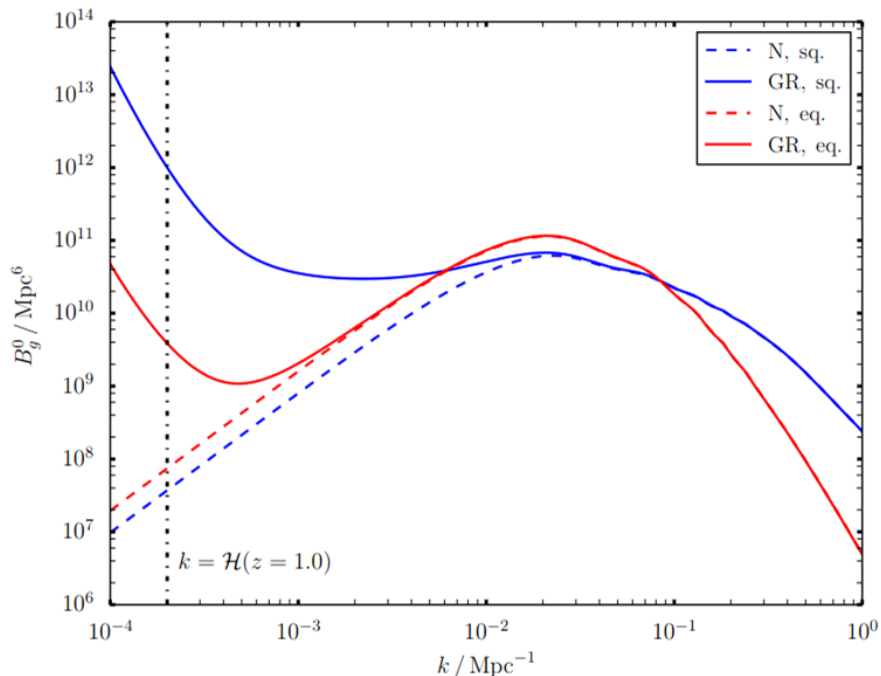
# Relativistic Galaxy Bispectrum



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$$B_g(\mathbf{k}_{123}) = \mathcal{K}^{(1)}(\mathbf{k}_1) \mathcal{K}^{(1)}(\mathbf{k}_2) \mathcal{K}^{(2)}(\mathbf{k}_{123}) P(\mathbf{k}_1) P(\mathbf{k}_2) + 2 \text{ c.c.}$$

$$\mathcal{K}^{(i)} = \mathcal{K}_N^{(i)} + \mathcal{K}_{GR}^{(i)}$$



Monopole of the bispectrum for squeezed and equilateral triangles

Shows a significant difference at larger scales i.e small  $k$

# Galaxy Surveys



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**Roman** -  $H\alpha$  spectroscopic survey from the Wide Field Instrument (WFI)

**DESI** - Bright Galaxy Sample (BGS)

**Euclid** -  $H\alpha$  sample from the Near Infrared Spectroscopic and Photometric (NISIP) instrument

**MegaMapper** - Future ground based wide field spectroscopic instrument

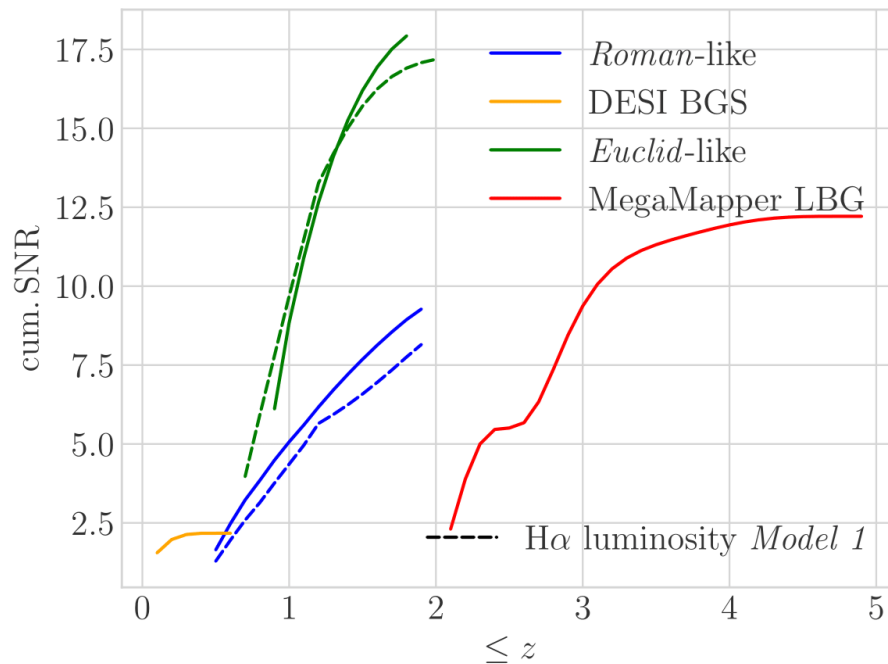
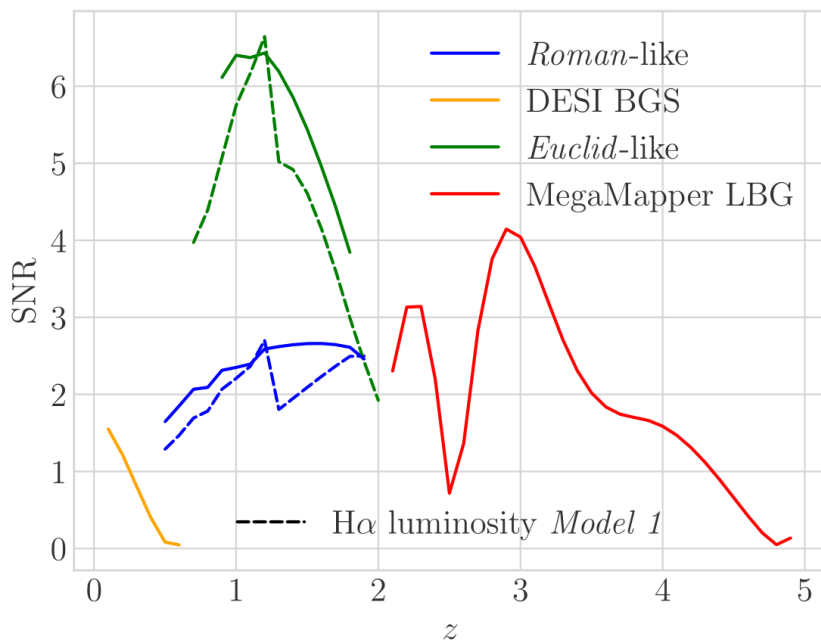
Survey	Redshift range	Sky area [deg <sup>2</sup> ]
Roman	<b>model 1 &amp; 3:</b> $0.5 \leq z \leq 2$	2000
DESI	$0 \leq z \leq 0.7$	14 000
Euclid	<b>model 1:</b> $0.7 \leq z \leq 2$ <b>model 3:</b> $0.9 \leq z \leq 1.8$	15 000
MegaMapper	$2 \leq z \leq 5$	20 000

# Signal-to-noise Ratio



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$$\left(\frac{S}{N}\right)_{GR}^2 = \sum_{k_a, \mu_1, \varphi} \frac{(B_g - B_N)(B_g - B_N)^*}{\text{Var}[B_{gN}]}$$



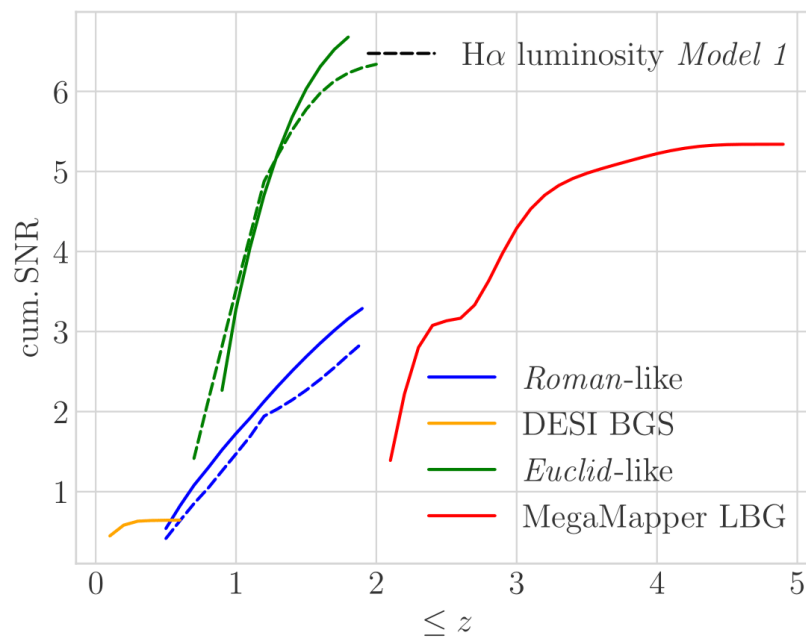
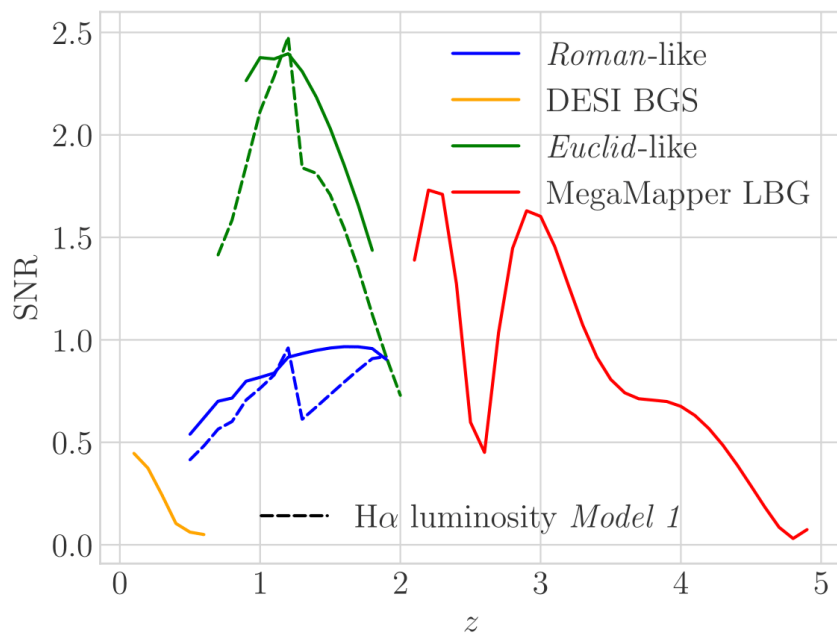


# Signal-to-noise Ratio



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$$\left(\frac{S}{N}\right)_{GR^{(2)}}^2 = \sum_{k_a, \mu_1, \varphi} \frac{(B_g - (B_N + B_{GR^{(1)}})) (B_g - (B_N + B_{GR^{(1)}}))^*}{\text{Var}[B_{gN}]}$$



# Marginal errors



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- We want to see how precisely we can measure the local PNG and relativistic contributions

$$\mathcal{K}^{(i)} = \mathcal{K}_N^{(i)} + \mathcal{K}_{\text{GR}}^{(i)} + \mathcal{K}_{\text{nG}}^{(i)}$$

**Fisher matrix formalism:**

$$F_{\alpha\beta} = \sum_{z, k_a, \mu_a, \varphi} \frac{\partial_{(\alpha} B_g \partial_{\beta)} B_g^*}{\text{Var}[B_g, B_g]}$$

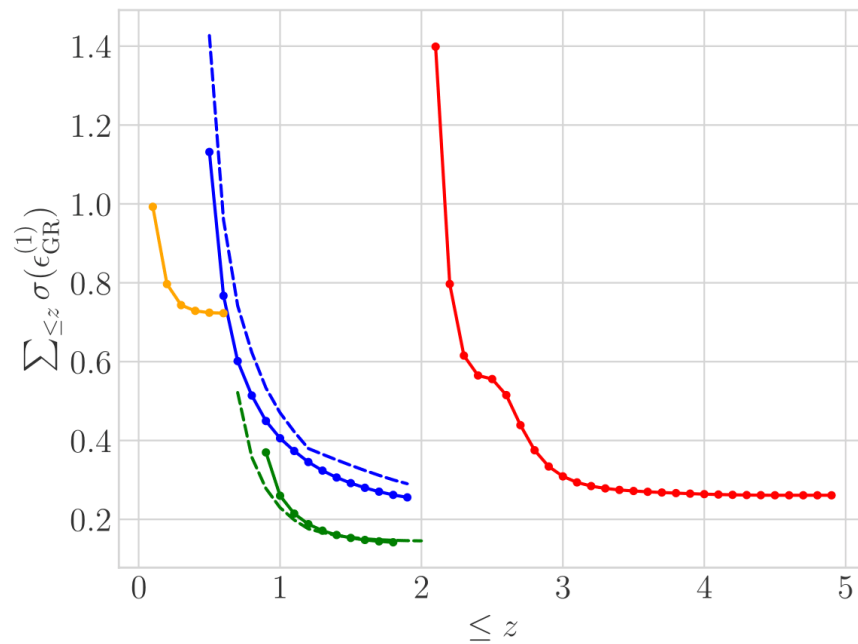
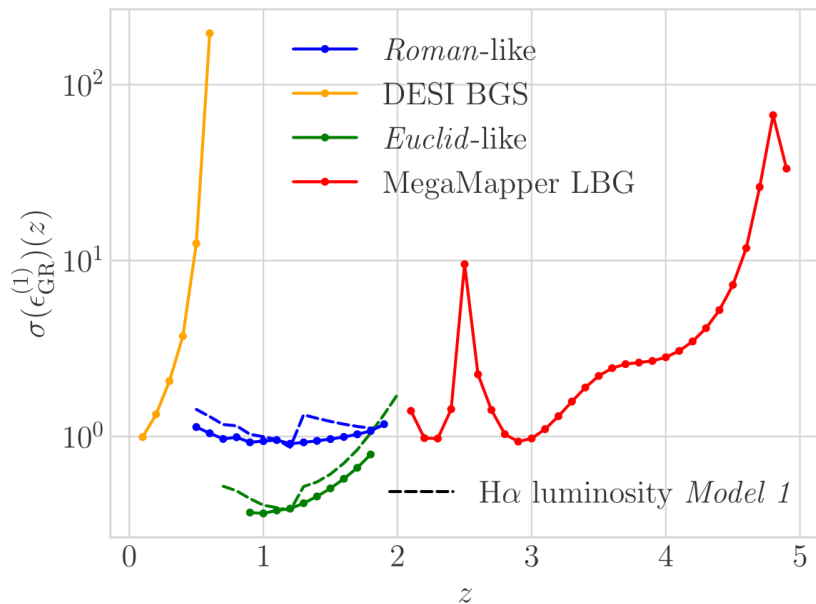
- From this we can get the marginal errors on our parameters  $\theta = \{\epsilon_{\text{GR}}^{(1)}, \epsilon_{\text{GR}}^{(2)}, f_{\text{NL}}\}$

$$\sigma_\alpha = \sqrt{(F^{-1})_{\alpha\alpha}}$$

# Marginal errors on $\epsilon_{GR}^{(1)}$



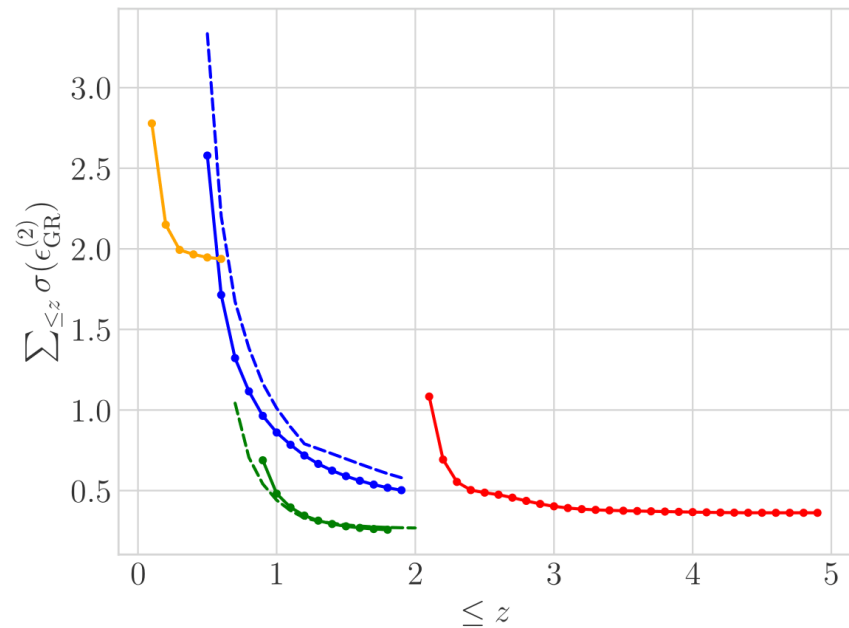
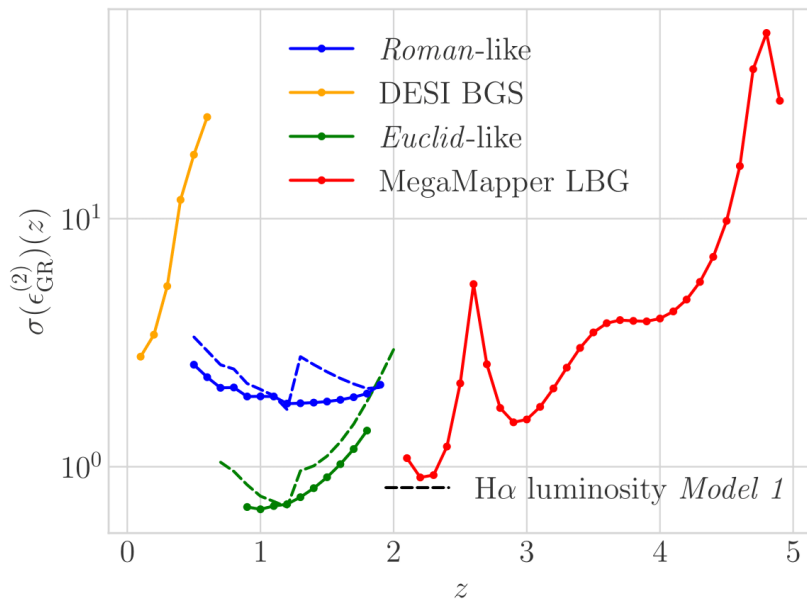
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# Marginal errors on $\epsilon_{GR}^{(2)}$



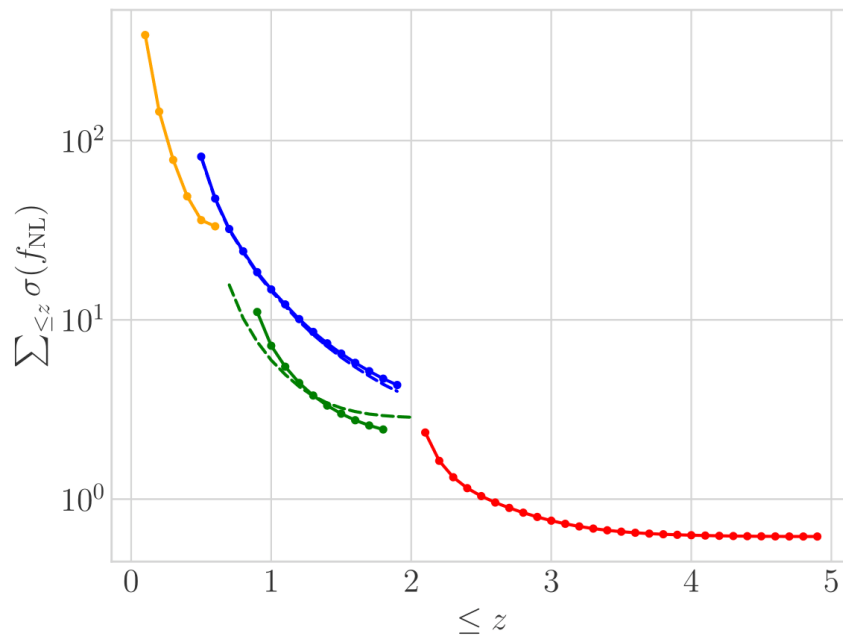
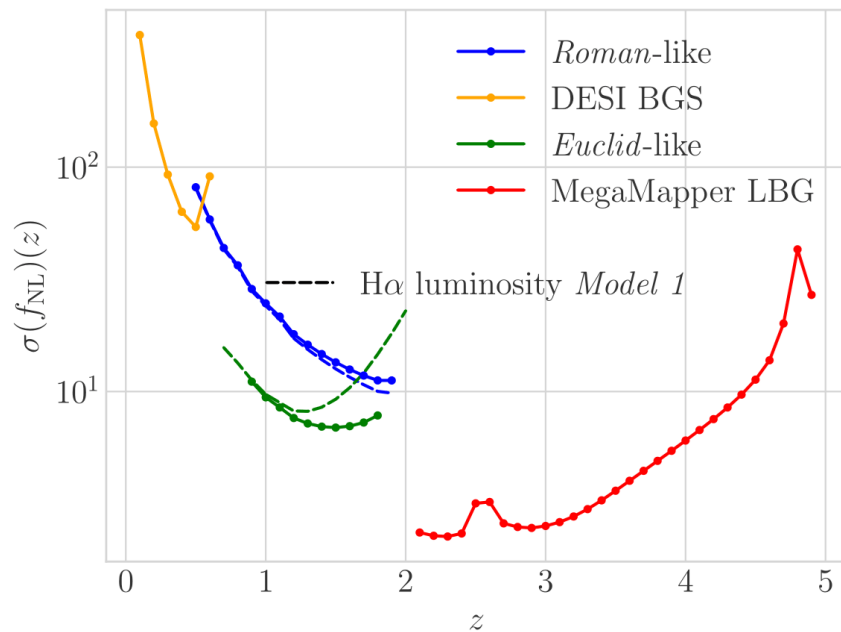
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# Marginal errors on $f_{\text{NL}}$



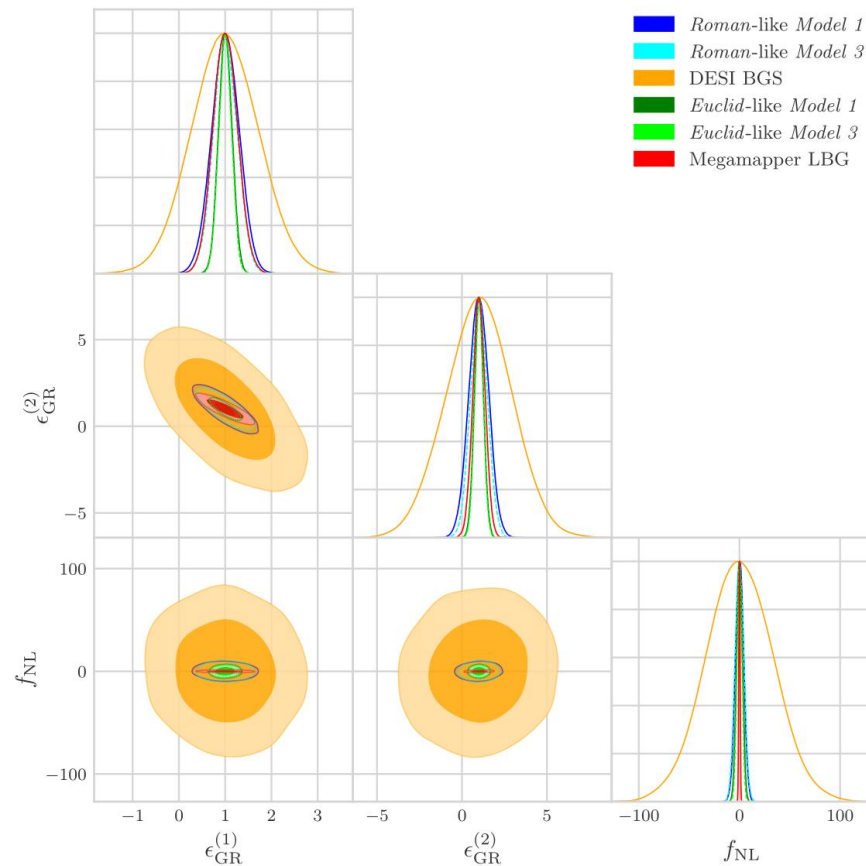
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# Constraints on all parameters



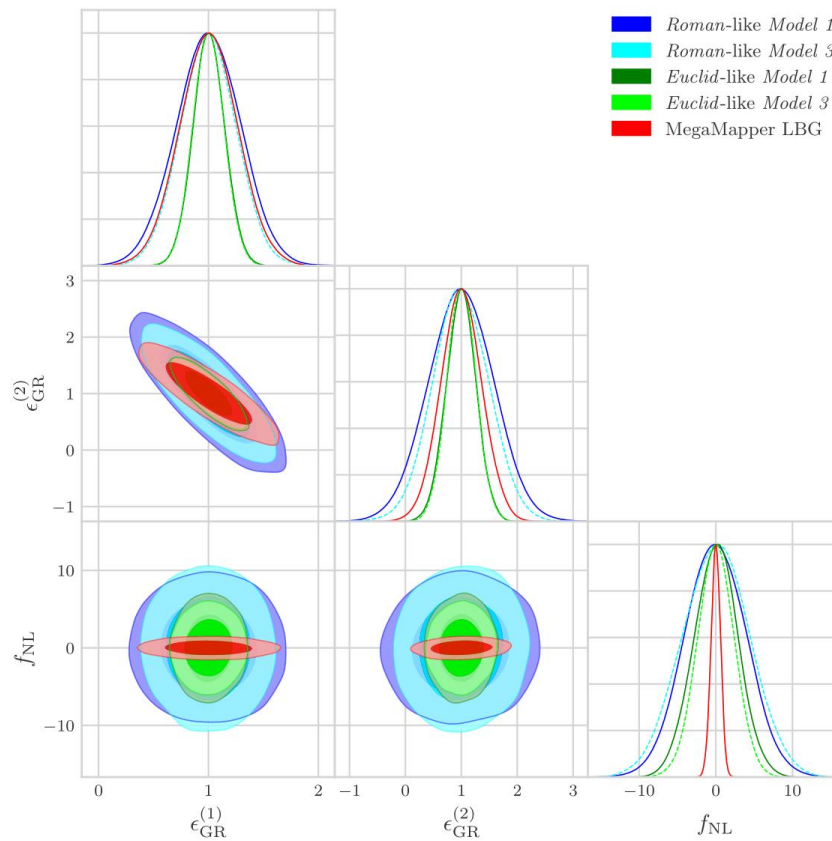
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# Constraints on all parameters



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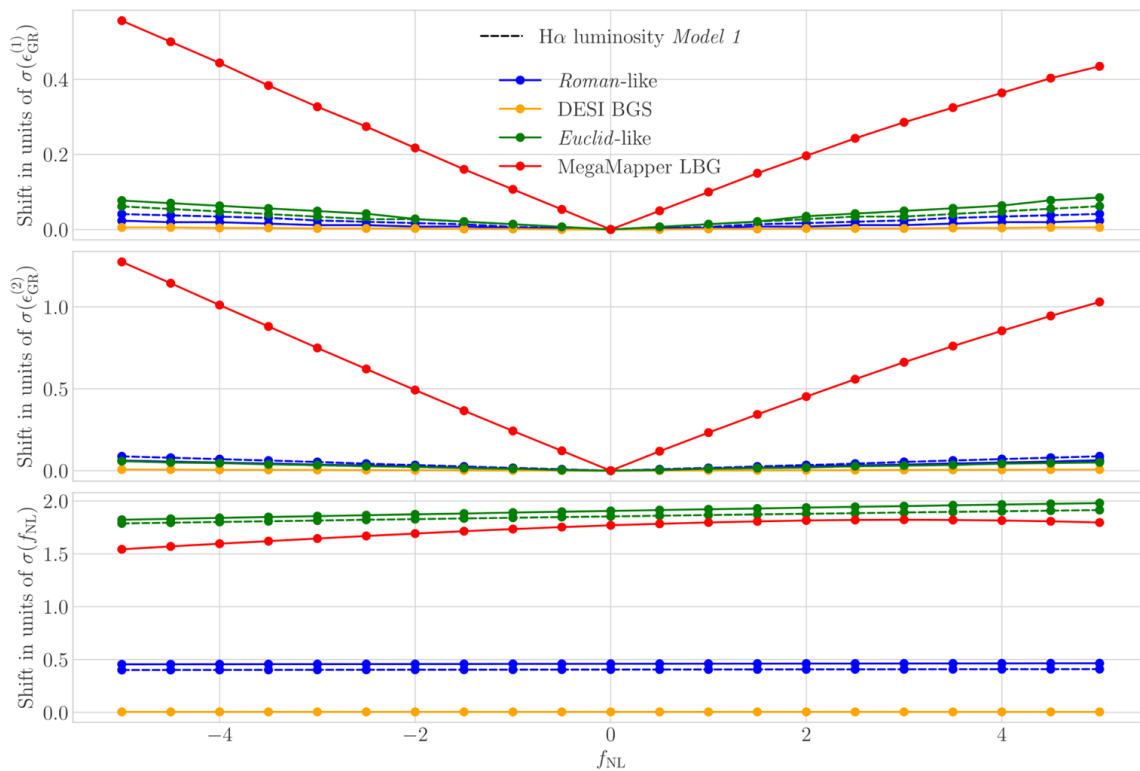


# Bias on parameters



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$$\Delta\theta_\alpha = \sum_\beta \left( F^{\theta_{GR}\theta_{GR}} \right)_{\alpha\beta}^{-1} F_\beta^{\theta f_{NL}} \delta f_{NL}$$



Uncertainty on  $f_{NL}$  makes little difference to the observed value we get on the relativistic contributions

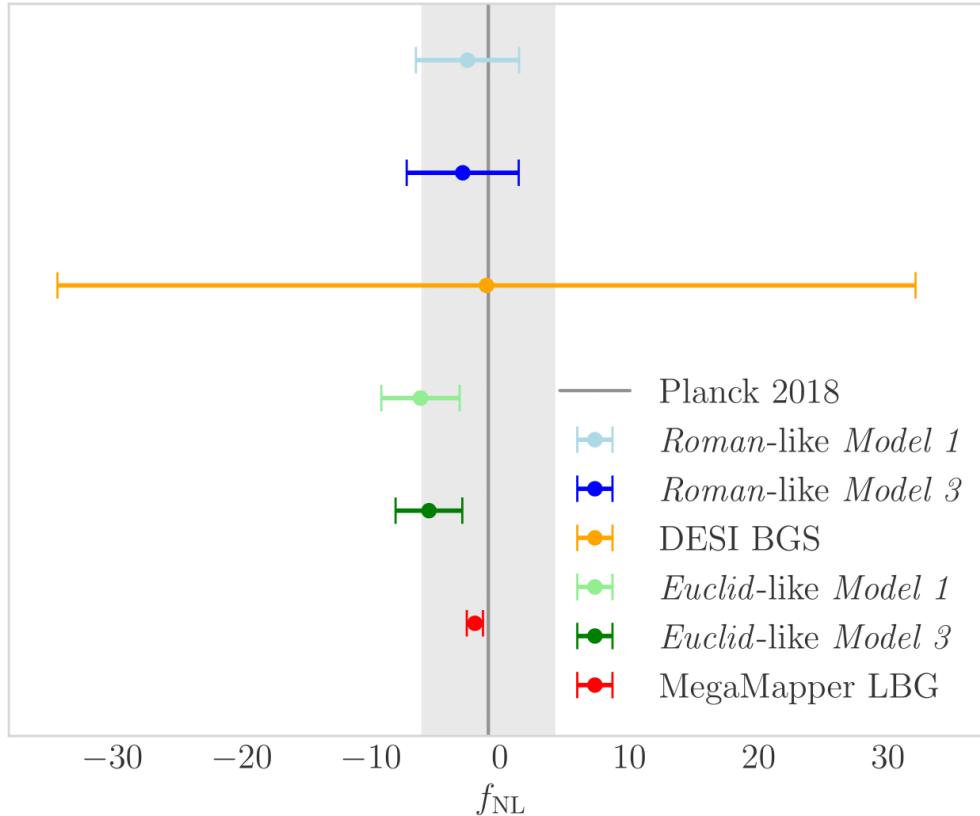
Neglecting Relativistic contributions would impact our observed value of  $f_{NL}$  significantly.



# Bias on parameters



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- The shift on the value of  $f_{\text{NL}}$  with error bars in the Newtonian regime.
- We have included the measured  $f_{\text{NL}}$  with error bars as from Planck 2018 as a reference (greyed out area).

# Cross Correlation Bispectrum $\implies$ Multi-tracer $\implies$ P + B



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$$B^{XXY}(\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3) = \frac{1}{3} \left\{ \mathcal{D}_B \left[ \left\{ \mathcal{K}_1^X(\mathbf{k}_1) \mathcal{K}_1^Y(\mathbf{k}_2) + \mathcal{K}_1^Y(\mathbf{k}_1) \mathcal{K}_1^X(\mathbf{k}_2) \right\} \mathcal{K}_2^X(\mathbf{k}_1, \mathbf{k}_2) \right. \right. \\ \left. \left. + \mathcal{K}_1^X(\mathbf{k}_1) \mathcal{K}_1^X(\mathbf{k}_2) \mathcal{K}_2^Y(\mathbf{k}_1, \mathbf{k}_2) \right] P(k_1) P(k_2) + 2\mathbf{k}\text{-perm} \right\},$$

$$B^{XYY}(\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3) = \frac{1}{3} \left\{ \mathcal{D}_B \left[ \left\{ \mathcal{K}_1^X(\mathbf{k}_1) \mathcal{K}_1^Y(\mathbf{k}_2) + \mathcal{K}_1^Y(\mathbf{k}_1) \mathcal{K}_1^X(\mathbf{k}_2) \right\} \mathcal{K}_2^Y(\mathbf{k}_1, \mathbf{k}_2) \right. \right. \\ \left. \left. + \mathcal{K}_1^Y(\mathbf{k}_1) \mathcal{K}_1^Y(\mathbf{k}_2) \mathcal{K}_2^X(\mathbf{k}_1, \mathbf{k}_2) \right] P(k_1) P(k_2) + 2\mathbf{k}\text{-perm} \right\},$$

Multi-tracer power spectra and  
bispectra: Formalism

[arXiv:2305.04028](https://arxiv.org/abs/2305.04028)

Dionysios Karagiannis,<sup>1</sup> Roy Maartens,<sup>1,2,3</sup> José Fonseca,<sup>4,5,1</sup>  
Stefano Camera,<sup>6,7,8,1</sup> Chris Clarkson<sup>9,1</sup>

# Summary



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- Relativistic effects couple with local  $f_{\text{NL}}^{\text{MC}}$  on large scales
- Without including relativistic effects,  $f_{\text{NL}}$  cannot be measured accurately
- The evolution bias leaves a big imprint on the detection of relativistic effects
- The modelling of the luminosity functions of surveys is crucial for decoupling
- Working with cross-correlations of different samples of the same survey may allow us to reduce the degeneracies within even the least promising surveys



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Thank you.  
Any questions?



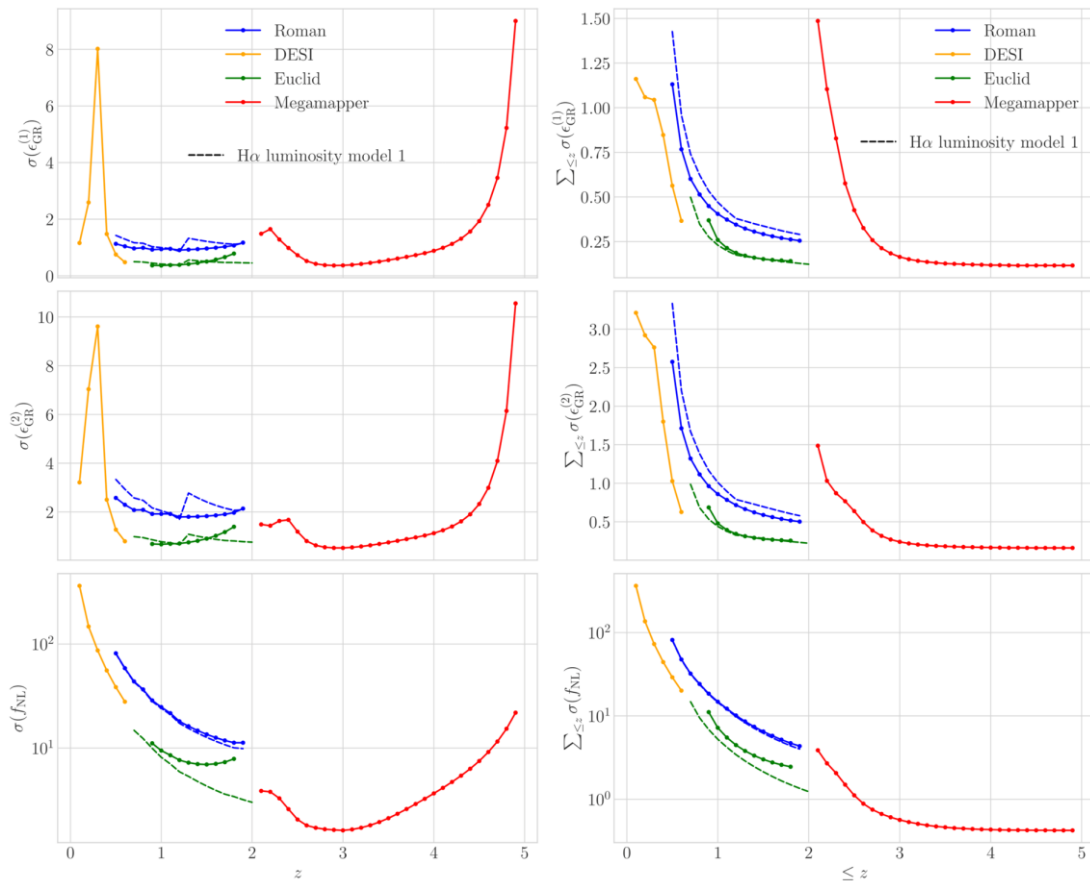
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Backup slides

# Marginal errors



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Survey	$\sigma(\epsilon_{\text{GR}}^{(1)})$	$\sigma(\epsilon_{\text{GR}}^{(2)})$	$\sigma(f_{\text{NL}})$
Roman model 1	0.290	0.579	4.007
Roman model 3	0.255	0.502	4.344
DESI	0.366	0.627	20.123
Euclid model 1	0.123	0.233	1.234
Euclid model 3	0.141	0.256	2.458
MegaMapper	0.116	0.161	0.424

# Signal-to-noise Redone



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- Discontinuity in luminosity function causes steep drop in signal at  $z \sim 1.3$
- Check against an improved luminosity model with updated Euclid specs

