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Relativistic matter bispectrum of cosmic structure



Thomas Montandon

Montandon et al (2404.02783) Montandon et al (2212.06799) Adamek et al (2110.11249)



Julian Adamek





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European Research Council

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• Why?

- Why?
 - LCDM says it
 - it is annoying for inflation

 $B(t,k_1,k_2,k_3) = T^{(1)}(k_1)T^{(1)}(k_2)T^{(1)}(k_3)B_{\mathcal{R}}(k_1,k_2,k_3) + 2T^{(1)}(k_1)T^{(1)}(k_2)T^{(2)}(k_1,k_2,k)P_{\mathcal{R}}(k_1)P_{\mathcal{R}}(k_2) + \text{perm.}$

GR and Radiation effects are degenerate in time and in momentum space with PNG!



 $+ 2\partial_{\theta} \left[-\psi_s^I - \frac{1}{H_s} (\partial_{\eta} \psi_s^I + \partial_{\tau} \psi_{0s}) \right] \int_{\eta_s}^{\eta_s} d\eta' \gamma_0^{ab} \partial_{\theta} \int_{\eta'}^{\eta_s} d\eta' \varphi^I (\eta'') \\ - 2a \psi_{\perp,s}^{a} \partial_{\theta} \int_{\eta_s}^{\eta_s} d\eta' \psi^I (\eta') + \frac{4}{r_s} \int_{\eta_s}^{\eta_s} d\eta' \left[\psi^I (\eta') \left(-\psi^I (\eta') - 2 \int_{\eta'}^{\eta_s} d\eta'' \partial_{\eta'} \psi^I (\eta'') \right) \right. \\ \left. + \gamma_0^{ab} \partial_{\theta} \left(\int_{\eta'}^{\eta_s} d\eta' \psi^I (\eta') \right) \partial_{\theta} \left(\int_{\eta'}^{\eta'} d\eta' \psi^I (\eta'') \right) \right] + 4 \partial_{\theta} \psi_s^I \int_{\eta_s}^{\eta_s} d\eta' \gamma_0^{ab} \partial_{\theta} \int_{\eta'}^{\eta_s} d\eta' \psi^I (\eta') \\ \left. - \frac{8}{2} \left[\int_{\eta'}^{\eta_s} d_{\eta'} \psi^I (\eta' \int_{\eta'}^{\eta_s} d\eta' \phi^I (\eta') \right] \right]$

 $\gamma_0^{ab} \partial_b \psi^I(\eta)$

 $\frac{1}{2}\Delta_2\psi^I(\eta')$

- Why?
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 - it is annoying
- How?
 - Computation

$$\begin{split} &+\frac{1}{\mathcal{H}_{s}}\partial_{\eta}\psi_{s}^{A}\left[-\left(\psi_{s}^{I}-\psi_{s}^{A}\right)-2\int_{\eta_{s}}^{\eta_{o}}d\eta'\psi^{I}(\eta')\right]+\frac{1}{\mathcal{H}_{s}}\psi_{s}^{A}\partial_{\eta}\psi_{s}^{J}-\\ &+\left[-\frac{1}{\mathcal{H}_{s}^{2}}\left(\partial_{\eta}^{2}\psi_{s}^{A}-\partial_{r}^{2}\psi_{s}^{A}-\partial_{\eta}\partial_{r}\psi_{s}^{A}\right)+\left(5+3\frac{\mathcal{H}_{s}'}{\mathcal{H}_{s}^{2}}+\frac{2}{\mathcal{H}_{s}r_{s}}\right)\frac{1}{\mathcal{H}_{s}}\partial_{\eta}\psi_{s}^{A}-\frac{3}{\mathcal{H}_{s}}\partial_{\tau}\psi_{s}^{A}\right.\\ &-\left(6+4\frac{\mathcal{H}_{s}'}{\mathcal{H}_{s}^{2}}\right)\psi_{s}^{A}\right]v_{||s}+\left[-\frac{1}{\mathcal{H}_{s}^{2}}\partial_{r}^{2}v_{||s}-\left(8+3\frac{\mathcal{H}_{s}'}{\mathcal{H}_{s}^{2}}\right)\frac{1}{\mathcal{H}_{s}}\partial_{\eta}v_{s}^{A}-\frac{2}{\mathcal{H}_{s}^{2}}\partial_{\eta}\psi_{s}^{A}\partial_{r}v_{||s}\right.\\ &+\frac{1}{\mathcal{H}_{s}}av_{\perp s}^{a}\partial_{a}\psi_{s}^{A}+\left[-\frac{2}{\mathcal{H}_{s}}\partial_{\eta}\psi_{s}^{A}+\left(10+2\frac{\mathcal{H}_{s}}{\mathcal{H}_{s}^{2}}\right)\psi_{s}^{A}\right]\frac{1}{r_{s}}\int_{\eta_{s}}^{\eta_{o}}d\eta'\frac{\eta'-\eta_{s}}{\eta_{o}-\eta'}\Delta_{2}\psi^{I}\left(\eta'\right)\\ &+\frac{2}{\mathcal{H}_{s}}\partial_{\tau}\psi_{s}^{A}\int^{\eta_{o}}d\eta'\Delta_{2}\psi^{I}\left(\eta'\right)+\left\{\left[-\frac{2}{\mathcal{H}_{s}}\partial_{\eta}\psi_{s}^{A}+2\left(4+\frac{\mathcal{H}_{s}'}{\mathcal{H}_{s}}\right)\psi_{s}^{A}\right]\frac{1}{1}\right] \\ \end{split}$$

$$\begin{split} &+\frac{4}{H_{h}r_{e}}\psi_{i}^{d}-\frac{2}{H_{h}r_{e}}\int_{0}^{\infty}d\theta A_{2}\psi_{i}^{d}\left(\eta\right)\left|\eta_{h}+\left[2\left(2+\frac{H_{h}^{d}}{H_{h}r_{e}}\right)\partial_{i}\eta_{h}\right)\partial_{i}\eta_{h}\right.\\ &+\frac{2}{H_{h}}\partial_{i}\eta_{h}^{d}\right]\int_{0}^{\infty}d\theta \psi_{i}^{d}\left(\eta\right)+\left[\frac{2}{H_{h}}\left(5+3\frac{H_{h}^{2}}{H_{h}^{2}}\right)\partial_{i}\eta_{h}+\frac{2}{H_{h}^{2}}\partial_{i}\eta_{h}\right)\int_{0}^{\infty}d\theta \partial_{i}\psi_{i}^{d}\left(\eta\right)\\ &+\left(6+3\frac{H_{h}^{2}}{H_{h}^{2}}\partial_{i}\eta_{h}-\frac{H_{h}^{2}}{\eta_{h}}\partial_{i}\eta_{h}^{d}\psi_{i}^{d}\left(\eta\right)+\frac{1}{H_{h}^{2}}\partial_{i}\psi_{h}^{d}\partial_{i}\psi_{i}^{d}\left(\eta\right)\\ &+\left(6+3\frac{H_{h}^{2}}{H_{h}^{2}}\partial_{i}\eta_{h}^{d}-\frac{H_{h}^{2}}{\eta_{h}}\partial_{i}\psi_{i}^{d}\eta_{h}^{d}\psi_{i}^{d}\left(\eta\right)\right)^{2}+\left\{\left[2\left(-2-\frac{H_{h}^{2}}{H_{h}^{2}}\right)\phi_{i}^{d}\right]\\ &+\left(6-3\frac{H_{h}^{2}}{H_{h}^{2}}\partial_{i}\phi_{i}^{d}\psi_{i}^{d}\left(\eta\right)-\frac{2}{H_{h}}\partial_{i}\phi_{i}^{d}\left(\eta\right)^{2}\right)^{2}+\left\{\left[2\left(-2-\frac{H_{h}^{2}}{H_{h}^{2}}\right)\phi_{i}^{d}\right]\\ &+4\left(-2-\frac{H_{h}^{2}}{H_{h}^{2}}\right)\int_{0}^{\infty}d\eta_{i}^{d}\psi_{i}^{d}\left(\eta\right)-\frac{2}{H_{h}}\partial_{i}\psi_{i}^{d}\left(\frac{1}{r_{h}}+2\left(-2-\frac{H_{h}^{2}}{H_{h}^{2}}\right)\partial_{i}\phi_{i}^{d}\partial_{i}\phi_{i}^{d}\left(\eta\right)\right)\\ &+2\left(-2-\frac{H_{h}^{2}}{H_{h}^{2}}\right)\partial_{i}\psi_{i}^{d}\left(-1-\frac{H_{h}^{2}}{H_{h}^{2}}\right)\partial_{i}\psi_{i}^{d}-\frac{1}{H_{h}^{2}}\partial_{i}\phi_{i}\psi_{i}^{d}\left(-2\int_{0}^{\infty}d\eta_{i}^{d}\psi_{i}^{d}\left(\eta\right)\right)\\ &+\left[\frac{3}{H_{h}^{2}}\left(-1-\frac{H_{h}^{2}}{H_{h}^{2}}-\frac{2}{3}\frac{1}{H_{h}r_{h}}\right)\partial_{i}\psi_{i}^{d}+\frac{1}{H_{h}^{2}}\partial_{i}\psi_{i}^{d}-\left(-2-\frac{H_{h}^{2}}{H_{h}^{2}}\right)\psi_{i}^{d}\right)\right. \end{aligned}$$

 $+\left(2-3\frac{\mathcal{H}_{s}'}{\mathcal{H}_{s}'}-\frac{2}{\mathcal{H}_{s}r_{s}}\right)\frac{1}{\mathcal{H}_{s}}\partial_{0}\psi_{s}^{\mathcal{A}}+\left(4-3\frac{\mathcal{H}_{s}'}{\mathcal{H}_{s}^{2}}-\frac{2}{\mathcal{H}_{s}r_{s}}\right)\frac{1}{\mathcal{H}_{s}}\partial_{0}\psi_{s}^{\mathcal{I}}-\left(7+\frac{\mathcal{H}_{s}'}{\mathcal{H}_{s}^{2}}\right)\psi_{s}^{\mathcal{I}}$

$$\begin{split} &+6\,\partial_a\psi_s^i\int_{\eta_c}^{\eta_c}d\eta'\gamma_0^{ab}\partial_b\int_{\eta'}^{\eta_c}d\eta'\eta_c^{ab}\langle\eta'^{-1}(\eta')+\frac{2}{\mathcal{H}_s}\partial_a\psi_s^i\gamma_0^{ab}\partial_b\int_{\eta_c}^{\eta_c}d\eta'\psi^I(\eta')\\ &-\frac{2}{\mathcal{H}_s}\partial_a\left(\partial_i\psi_s^i\right)\int_{\eta_c}^{\eta_c}d\eta'\gamma_0^{ab}\partial_b\int_{\eta'}^{\eta_c}d\eta'\psi^I(\eta')+\frac{8}{r_s}\int_{\eta_c}^{\eta_c}d\eta'\left(\psi^I\psi^A\right)(\eta')\\ &-\frac{4}{r_s}\int_{\eta_c}^{\eta_c}d\eta'\frac{\eta'-\eta_s}{\eta_c-\eta'}\Delta_2\left(\psi^I\psi^A\right)(\eta')+\left[\left(-\frac{2}{\mathcal{H}_sr_s}-\frac{\mathcal{H}_s'}{\mathcal{H}_s^2}\right)\psi_s^i+\frac{1}{\mathcal{H}_s}\partial_\eta\psi_s^A\right]\delta_{\mu}^{[1)}\\ &+\left[\frac{1}{\rho}\partial_q\left(\bar{\rho}\,\delta_{\mu}^{(1)}\right)-\partial_s\delta_{\mu}^{(1)}\right]\frac{1}{\mathcal{H}_s}\psi_s^i. \end{split} \tag{4}$$

- Why?
 - LCDM says it
 - it is annoying for inflation
- How?
 - Computation is too hard...
 - SIMULATION!



Inflaton ϕ



Inflaton ϕ













Results







- Redshift bins
- GR up to \mathcal{H}^4/k^4
- Radiation up to \mathcal{H}^4/k^4
- Projection effects up to \mathcal{H}/k

 $egin{aligned} b_{\ell_1\ell_2\ell_3}^{r_1r_2r_3} &= rac{8}{\pi^3}\int dr_1' dr_2' dr_3' W(r_1,r_1') W(r_2,r_2') W(r_3,r_3') \ \int dk_1 dk_2 dk_3 d\chi (k_1k_2k_3\chi)^2 B^{r_1'r_2'r_3'}(k_1,k_2,k_3) \ j_{\ell_1}(k_1r_1') j_{\ell_2}(k_2r_2') j_{\ell_3}(k_3r_3') j_{\ell_1}(k_1\chi) j_{\ell_2}(k_2\chi) j_{\ell_3}(k_3\chi) \end{aligned}$

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$$b_{\ell_{1}\ell_{2}\ell_{3}}^{\hat{z}_{1}\hat{z}_{2}\hat{z}_{3}} = b_{\ell_{1}\ell_{2}\ell_{3}}^{\delta_{2}} + b_{\ell_{1}\ell_{2}\ell_{3}}^{\partial_{r}^{2}v_{1}} + b_{\ell_{1}\ell_{2}\ell_{3}}^{\partial_{r}v_{1}\partial_{r}^{3}v_{1}} + b_{\ell_{1}\ell_{2}\ell_{3}}^{\partial_{r}v_{1}\partial_{r}\delta_{1}} + b_{\ell_{1}\ell_{2}\ell_{3}}^{\partial_{r}v_{1}\partial$$

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$$b_{\ell_{1}\ell_{2}\ell_{3}}^{\hat{z}_{1}\hat{z}_{2}\hat{z}_{3}} = b_{\ell_{1}\ell_{2}\ell_{3}}^{\delta_{2}} + b_{\ell_{1}\ell_{2}\ell_{3}}^{\partial_{r}^{2}v_{1}} + b_{\ell_{1}\ell_{2}\ell_{3}}^{\partial_{r}v_{1}\partial_{r}^{3}v_{1}} + b_{\ell_{1}\ell_{2}\ell_{3}}^{\partial_{r}v_{1}\partial_{r}\delta_{1}} + b_{\ell_{1}\ell_{2}\ell_{3}}^{\partial_{r}v_{1}\partial$$

$$b_{\ell_{1}\ell_{2}\ell_{3}}^{\delta_{2}} = 2 \sum_{mn} \int d\chi C_{\ell_{2}}^{(n,0,0)}(\chi) C_{\ell_{3}}^{(m,0,0)}(\chi) \left(\chi^{2} \int \frac{dr_{1}}{r_{1}^{2}} D^{2}(r_{1}) W(r_{1}) \left[f_{nm}^{(-4)} I_{\ell_{1}}(-1,r_{1},\chi) + f_{nm}^{(-2)} A_{\ell_{1}}(r_{1},\chi) \right. \left. + \sum_{p} c_{p} \left(f_{nm}^{(-4,R)} I_{\ell_{1}}(b+ip\eta-1,r_{1},\chi) + f_{nm}^{(-2,R)} I_{\ell_{1}}(b+ip\eta+1,r_{1},\chi) \right) \right] \left. + f_{nm}^{(0)} D^{2}(\chi) W(\chi) + \mathcal{D}_{\ell_{1}} \left[f_{nm}^{(2)} D^{2}(\chi) W(\chi) \right] + \mathcal{D}_{\ell_{1}}^{2} \left[f_{nm}^{(4)} D^{2}(\chi) W(\chi) \right] \right).$$
(3.14)



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- Redshift bins
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$$b_{ ext{sim}}^{ ext{GR}} \stackrel{?}{=} b_{ ext{th}}^{ ext{GR}} \quad ext{and} \quad b_{ ext{sim}}^{ ext{Rad}} \stackrel{?}{=} b_{ ext{th}}^{ ext{Rad}}$$

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$$b_{ ext{sim}}^{ ext{GR}} \stackrel{?}{=} b_{ ext{th}}^{ ext{GR}} \quad ext{and} \quad b_{ ext{sim}}^{ ext{Rad}} \stackrel{?}{=} b_{ ext{th}}^{ ext{Rad}}$$

Come back soon for the results, we are on it



GR effects are cool

- We have an accurate numerical model thanks to simulation
 - Theory is hard as expected

Thank you !



