



# The peculiar velocity bispectrum

*Stefano Camera*

Department of Physics, Alma Felix University of Turin, Italy



Funded by  
the European Union  
NextGenerationEU

UNIVERSITÀ  
DI TORINO



# Peculiar velocities



UNIVERSITÀ  
DI TORINO

- Galaxies' peculiar velocities can be estimated by combining measurements of their observed redshift with a distance indicator that enables us to infer the cosmological redshift independently

[Davis & Scrimgeour 2014; Watkins & Feldman 2015]

# Peculiar velocities

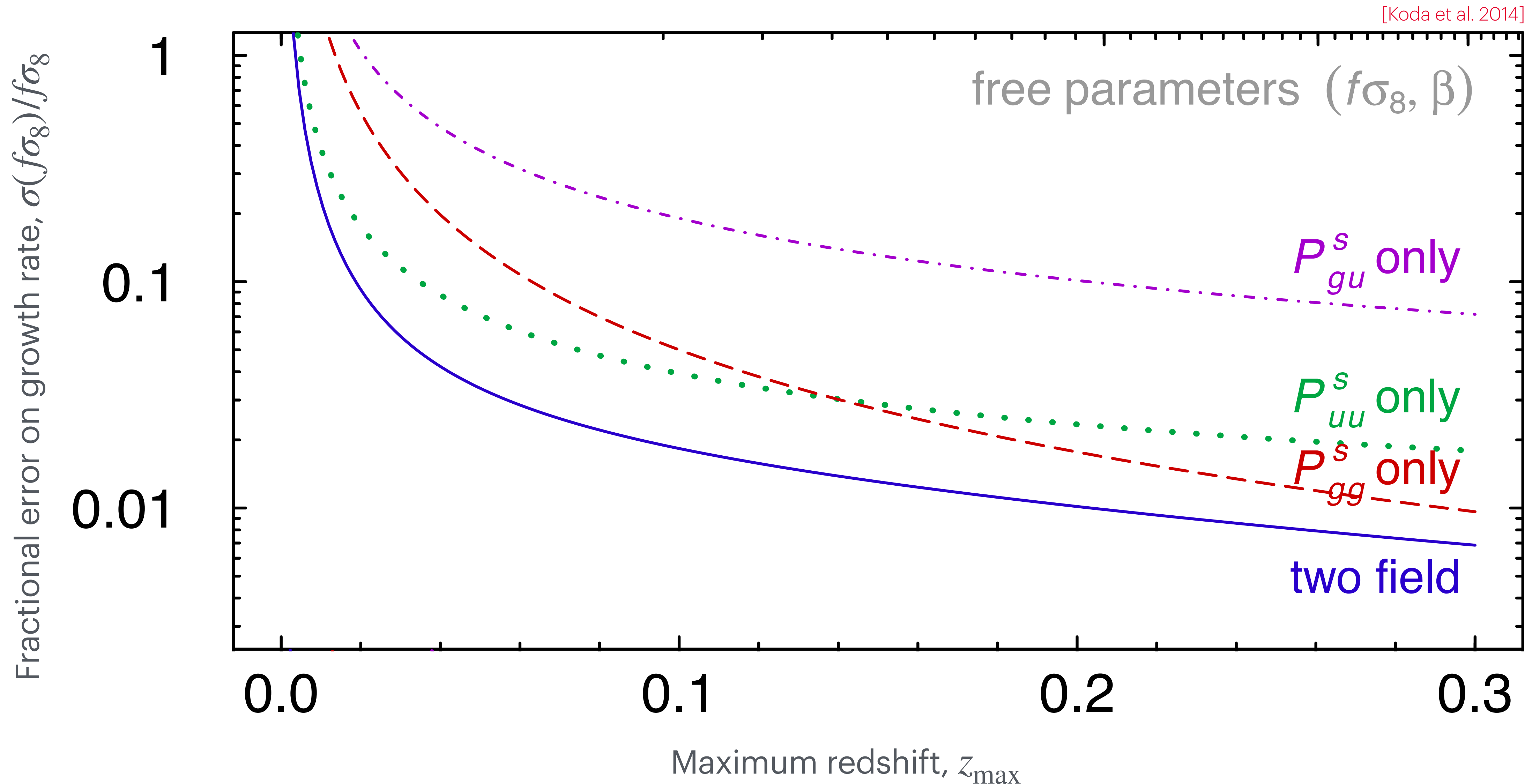


- Galaxies' peculiar velocities can be estimated by combining measurements of their observed redshift with a distance indicator that enables us to infer the cosmological redshift independently
- The power spectrum of peculiar velocities has been proposed as a powerful tool to complement traditional galaxy clustering

[Davis & Scrimgeour 2014; Watkins & Feldman 2015]

[Burkey & Taylor 2004; Iršič & Slosar 2011; Koda et al. 2014; Howlett et al. 2017a,b; Whitford et al. 2021]

# Peculiar velocities



# Peculiar velocities



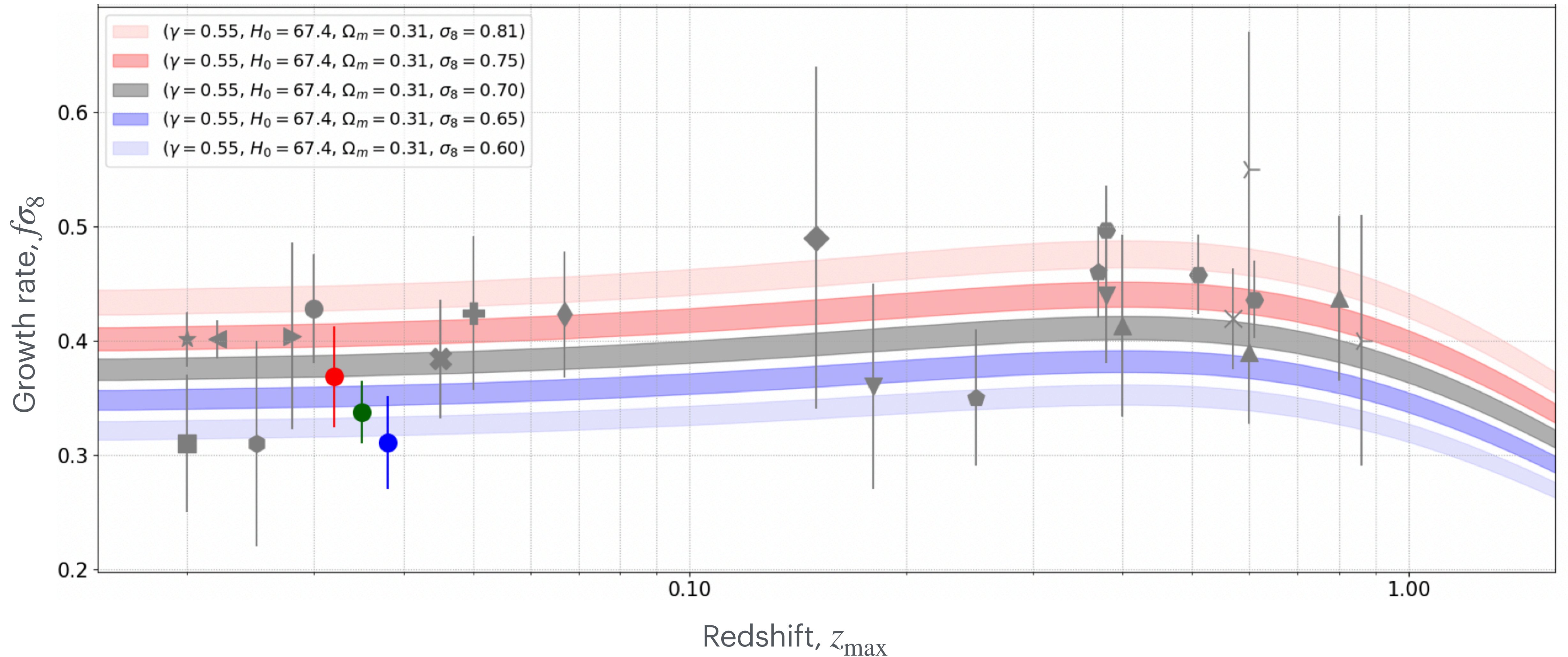
UNIVERSITÀ  
DI TORINO

- Galaxies' **peculiar velocities** can be estimated by combining measurements of their **observed redshift** with a distance indicator that enables us to infer the **cosmological redshift** independently  
[Davis & Scrimgeour 2014; Watkins & Feldman 2015]
- The **power spectrum** of peculiar velocities has been proposed as a powerful tool to complement traditional **galaxy clustering**  
[Burkey & Taylor 2004; Iršič & Slosar 2011; Koda et al. 2014; Howlett et al. 2017a,b; Whitford et al. 2021]
- Surveys that directly measure peculiar velocities have demonstrated dramatic improvement on constraints on the **growth of structure** relative to galaxy clustering alone at low redshift  
[Carrick et al. 2015; Qin et al. 2019; Adams & Blake 2020; Said et al. 2020; Lai et al. 2022]

# Peculiar velocities



[Said et al. 2020]



# Peculiar velocities



- Galaxies' peculiar velocities can be estimated by combining measurements of their observed redshift with a distance indicator that enables us to infer the cosmological redshift independently  
[Davis & Scrimgeour 2014; Watkins & Feldman 2015]
- The power spectrum of peculiar velocities has been proposed as a powerful tool to complement traditional galaxy clustering  
[Burkey & Taylor 2004; Iršič & Slosar 2011; Koda et al. 2014; Howlett et al. 2017a,b; Whitford et al. 2021]
- Surveys that directly measure peculiar velocities have demonstrated dramatic improvement on constraints on the growth of structure relative to galaxy clustering alone at low redshift  
[Carrick et al. 2015; Qin et al. 2019; Adams & Blake 2020; Said et al. 2020; Lai et al. 2022]
- What about the bispectrum, then?

# Polyspectra formalism



UNIVERSITÀ  
DI TORINO

- Given an observable **cosmological perturbation field**  $X$  (e.g. galaxy number counts, peculiar velocities, ...):



# Polyspectra formalism



- Given an observable **cosmological perturbation field**  $X$  (e.g. galaxy number counts, peculiar velocities, ...):
  - Fourier-space summary statistics for  $N$ -point correlation functions [**polyspectra**]
  - Auto-correlations only

# Polyspectra formalism



- Given an observable **cosmological perturbation field**  $X$  (e.g. galaxy number counts, peculiar velocities, ...):
  - Fourier-space summary statistics for  $N$ -point correlation functions [**polyspectra**]
  - Auto-correlations only
- **Power spectrum**

$$\langle X(\mathbf{k}_1) X(\mathbf{k}_2) \rangle = (2\pi)^3 \delta_{\text{D}}(\mathbf{k}_{12}) P_X(\mathbf{k}_1)$$

# Polyspectra formalism

- Given an observable **cosmological perturbation field**  $X$  (e.g. galaxy number counts, peculiar velocities, ...):
  - Fourier-space summary statistics for  $N$ -point correlation functions [**polyspectra**]
  - Auto-correlations only
- **Power spectrum**

$$\langle X(\mathbf{k}_1) X(\mathbf{k}_2) \rangle = (2\pi)^3 \delta_{\text{D}}(\mathbf{k}_{12}) P_X(\mathbf{k}_1)$$

- **Bispectrum**

$$\langle X(\mathbf{k}_1) X(\mathbf{k}_2) X(\mathbf{k}_3) \rangle = (2\pi)^3 \delta_{\text{D}}(\mathbf{k}_{123}) B_X(\mathbf{k}_1, \mathbf{k}_2)$$

# Standard perturbation theory primer



UNIVERSITÀ  
DI TORINO

- Small perturbations can be treated **perturbatively**

# Standard perturbation theory primer



UNIVERSITÀ  
DI TORINO

- Small perturbations can be treated **perturbatively**
- As long as the matter distribution can be described as an **irrotational fluid** (i.e. no shell crossing), gravitational instability can be fully described by the **density contrast  $\delta$**  and the **velocity divergence  $\theta$**

# Standard perturbation theory primer



UNIVERSITÀ  
DI TORINO

- Small perturbations can be treated **perturbatively**
- As long as the matter distribution can be described as an **irrotational fluid** (i.e. no shell crossing), gravitational instability can be fully described by the **density contrast  $\delta$**  and the **velocity divergence  $\theta$**
- In the **spherical collapse model** (sensible for small perturbations), solutions at any order remain **separable** in time and scale

# Standard perturbation theory primer



- Small perturbations can be treated **perturbatively**
- As long as the matter distribution can be described as an **irrotational fluid** (i.e. no shell crossing), gravitational instability can be fully described by the **density contrast**  $\delta$  and the **velocity divergence**  $\theta$
- In the **spherical collapse model** (sensible for small perturbations), solutions at any order remain **separable** in time and scale

$$\delta(\mathbf{k}, z) = \sum_{m=1}^{\infty} D^m(z) \delta^{(m)}(\mathbf{k})$$

$$\theta(\mathbf{k}, z) = -\mathcal{H}(z) f(z) \sum_{m=1}^{\infty} D^m(z) \theta^{(m)}(\mathbf{k})$$

# Standard perturbation theory primer



- From the continuity and Euler's equations in Fourier space

$$\begin{bmatrix} \delta^{(m)}(\mathbf{k}) \\ \theta^{(m)}(\mathbf{k}) \end{bmatrix} = \int_{\mathbf{q}_1} \cdots \int_{\mathbf{q}_m} \delta^{(1)}(\mathbf{q}_1) \cdots \delta^{(1)}(\mathbf{q}_m) \begin{bmatrix} F_m(\mathbf{q}_1, \dots, \mathbf{q}_m) \\ G_m(\mathbf{q}_1, \dots, \mathbf{q}_m) \end{bmatrix} (2\pi)^3 \delta_D(\mathbf{q}_{1\dots m} - \mathbf{k})$$



# Standard perturbation theory primer



- From the **continuity** and **Euler's equations** in Fourier space

$$\begin{bmatrix} \delta^{(m)}(\mathbf{k}) \\ \theta^{(m)}(\mathbf{k}) \end{bmatrix} = \int_{\mathbf{q}_1} \cdots \int_{\mathbf{q}_m} \delta^{(1)}(\mathbf{q}_1) \cdots \delta^{(1)}(\mathbf{q}_m) \begin{bmatrix} F_m(\mathbf{q}_1, \dots, \mathbf{q}_m) \\ G_m(\mathbf{q}_1, \dots, \mathbf{q}_m) \end{bmatrix} (2\pi)^3 \delta_D(\mathbf{q}_{1\dots m} - \mathbf{k})$$

- Hence the **tree-level** power spectrum and bispectrum

$$\langle X(\mathbf{k}_1) X(\mathbf{k}_2) \rangle = \langle X^{(1)}(\mathbf{k}_1) X^{(1)}(\mathbf{k}_2) \rangle$$

$$\langle X(\mathbf{k}_1) X(\mathbf{k}_2) X(\mathbf{k}_3) \rangle = \langle X^{(1)}(\mathbf{k}_1) X^{(1)}(\mathbf{k}_2) X^{(2)}(\mathbf{k}_3) \rangle + 2 \circlearrowleft \mathbf{k}_i$$

# Galaxy clustering



- Galaxy clustering is the summary statistics of **fluctuations in galaxy number counts**

$$\delta_g(\mathbf{x}) := \frac{n_g(\mathbf{x}) - \bar{n}_g}{\bar{n}_g}$$

# Galaxy clustering



- Galaxy clustering is the summary statistics of **fluctuations in galaxy number counts**

$$\delta_g(\mathbf{x}) := \frac{n_g(\mathbf{x}) - \bar{n}_g}{\bar{n}_g}$$

- **Redshift-space distortions (RSD)** arise as we don't know galaxies' distances, but infer them from their observed redshifts, which include **radial peculiar velocities**

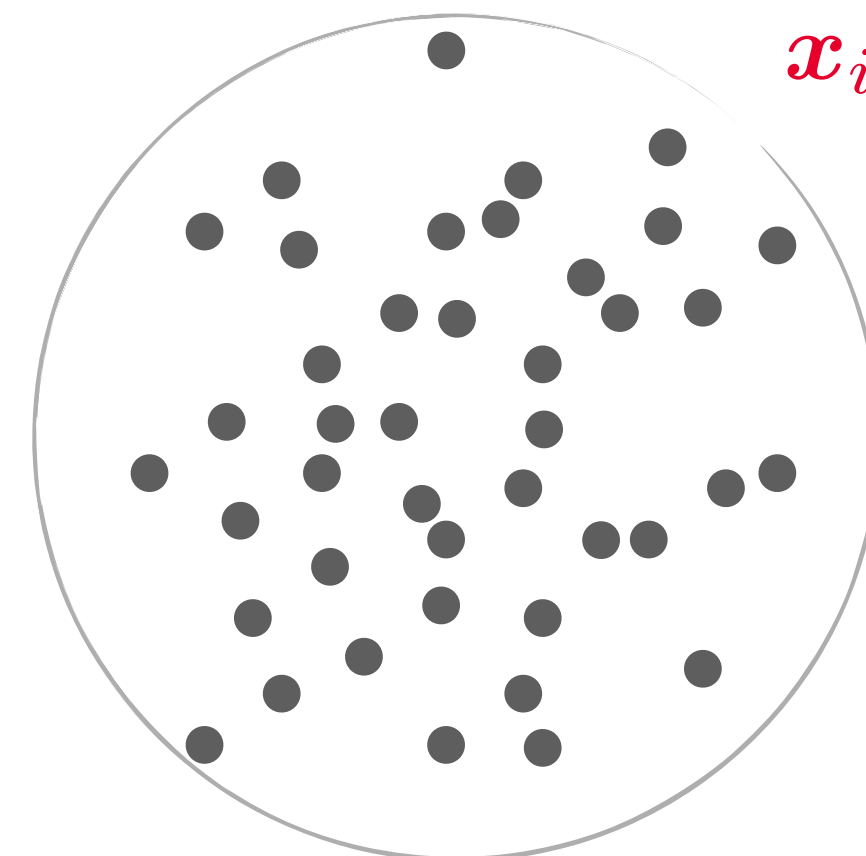
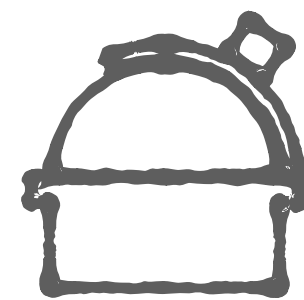
# Galaxy clustering



- Galaxy clustering is the summary statistics of **fluctuations in galaxy number counts**

$$\delta_g(\mathbf{x}) := \frac{n_g(\mathbf{x}) - \bar{n}_g}{\bar{n}_g}$$

- Redshift-space distortions (RSD)** arise as we don't know galaxies' distances, but infer them from their observed redshifts, which include **radial peculiar velocities**



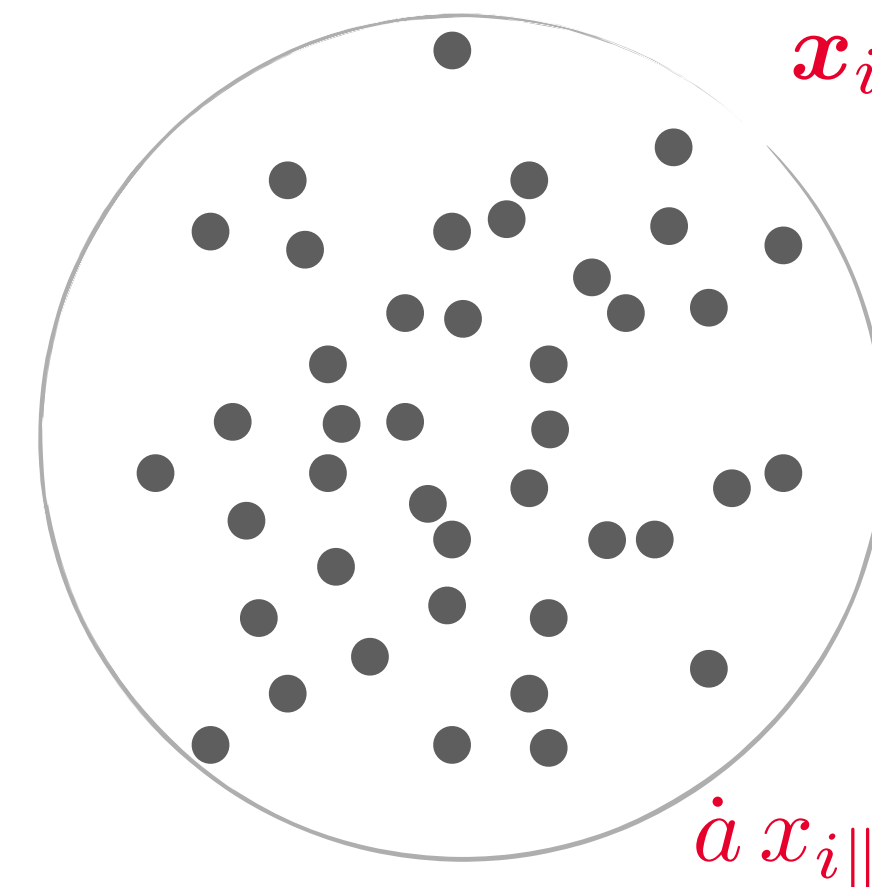
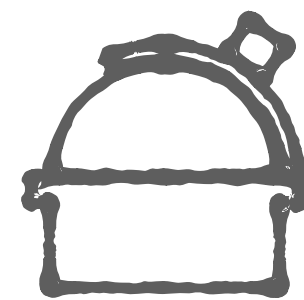
# Galaxy clustering



- Galaxy clustering is the summary statistics of **fluctuations in galaxy number counts**

$$\delta_g(\mathbf{x}) := \frac{n_g(\mathbf{x}) - \bar{n}_g}{\bar{n}_g}$$

- Redshift-space distortions (RSD)** arise as we don't know galaxies' distances, but infer them from their observed redshifts, which include **radial peculiar velocities**



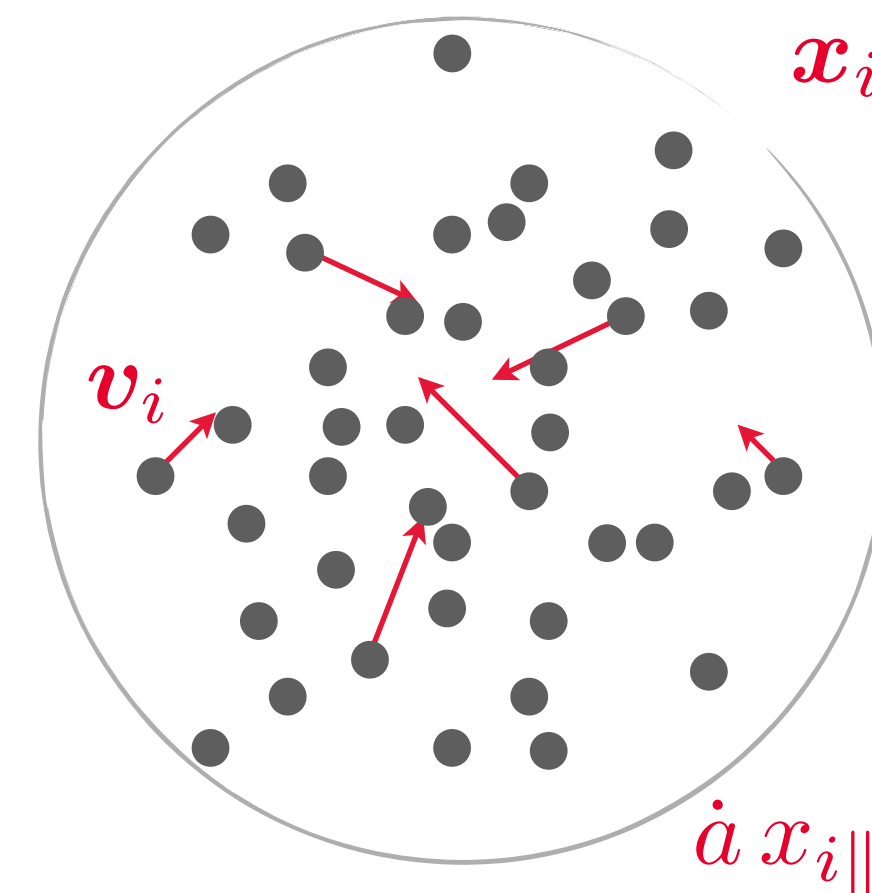
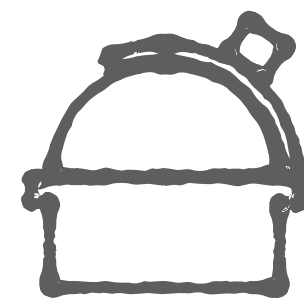
# Galaxy clustering



- Galaxy clustering is the summary statistics of **fluctuations in galaxy number counts**

$$\delta_g(\mathbf{x}) := \frac{n_g(\mathbf{x}) - \bar{n}_g}{\bar{n}_g}$$

- Redshift-space distortions (RSD)** arise as we don't know galaxies' distances, but infer them from their observed redshifts, which include **radial peculiar velocities**



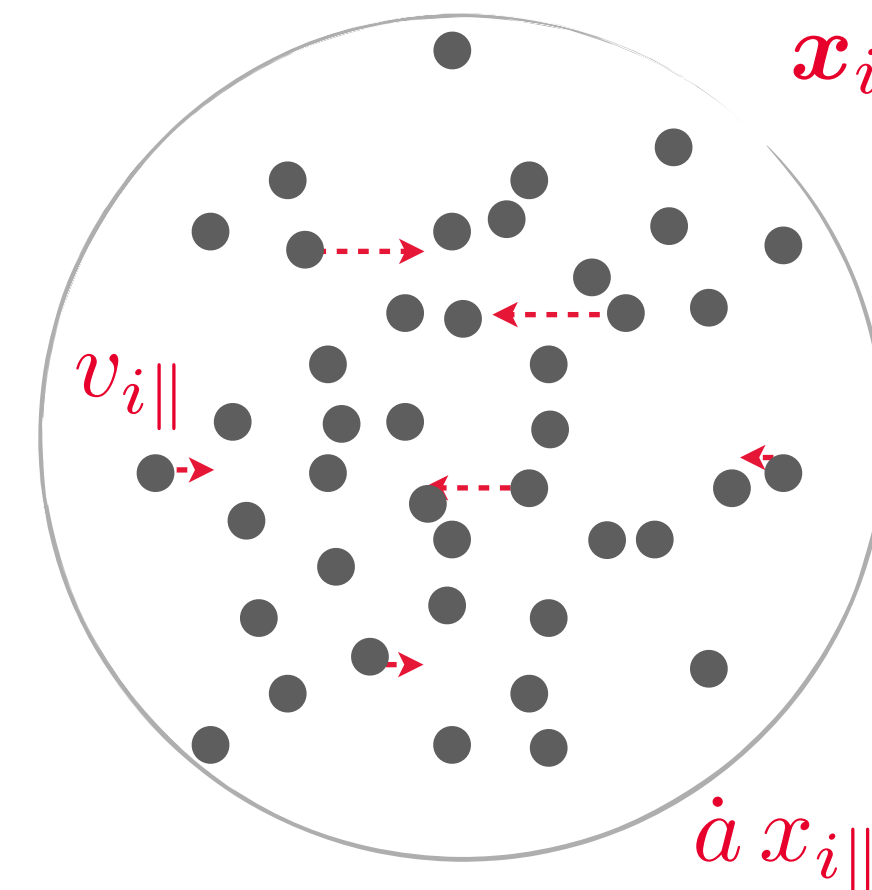
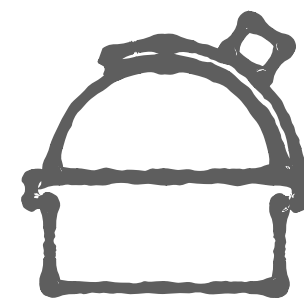
# Galaxy clustering



- Galaxy clustering is the summary statistics of **fluctuations in galaxy number counts**

$$\delta_g(\mathbf{x}) := \frac{n_g(\mathbf{x}) - \bar{n}_g}{\bar{n}_g}$$

- Redshift-space distortions (RSD)** arise as we don't know galaxies' distances, but infer them from their observed redshifts, which include **radial peculiar velocities**



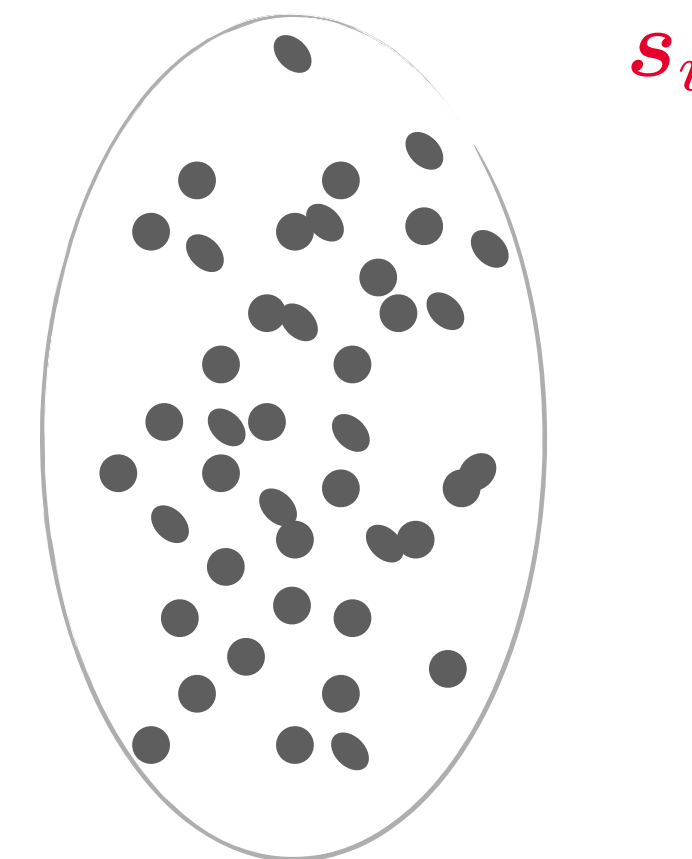
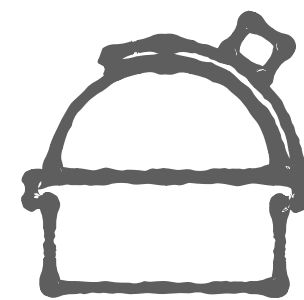
# Galaxy clustering



- Galaxy clustering is the summary statistics of **fluctuations in galaxy number counts**

$$\delta_g(\mathbf{x}) := \frac{n_g(\mathbf{x}) - \bar{n}_g}{\bar{n}_g}$$

- Redshift-space distortions (RSD)** arise as we don't know galaxies' distances, but infer them from their observed redshifts, which include **radial peculiar velocities**





# Galaxy clustering



UNIVERSITÀ  
DI TORINO

- However, number density conservation dictates that

$$d^3s [1 + \Delta(\mathbf{s})]$$

# Galaxy clustering



- However, number density conservation dictates that

$$d^3s [1 + \Delta(\mathbf{s})] = d^3x [1 + \delta_g(\mathbf{x})]$$

# Galaxy clustering



- However, **number density conservation** dictates that

$$d^3s [1 + \Delta(\mathbf{s})] = d^3x [1 + \delta_g(\mathbf{x})]$$

- Knowing the real-to-redshift space mapping, we construct **redshift-space kernels...**

$$\Delta^{(m)}(\mathbf{k}) = \int_{\mathbf{q}_1} \dots \int_{\mathbf{q}_m} \delta^{(1)}(\mathbf{q}_1) \dots \delta^{(1)}(\mathbf{q}_m) \mathcal{Z}^{(m)}(\mathbf{q}_1, \dots, \mathbf{q}_m) (2\pi)^3 \delta_D(\mathbf{q}_{1\dots m} - \mathbf{k})$$

# Galaxy clustering



- However, **number density conservation** dictates that

$$d^3s [1 + \Delta(\mathbf{s})] = d^3x [1 + \delta_g(\mathbf{x})]$$

- Knowing the real-to-redshift space mapping, we construct **redshift-space kernels**...

$$\Delta^{(m)}(\mathbf{k}) = \int_{\mathbf{q}_1} \dots \int_{\mathbf{q}_m} \delta^{(1)}(\mathbf{q}_1) \dots \delta^{(1)}(\mathbf{q}_m) \mathcal{Z}^{(m)}(\mathbf{q}_1, \dots, \mathbf{q}_m) (2\pi)^3 \delta_D(\mathbf{q}_{1\dots m} - \mathbf{k})$$

- ...from which

$$P_\Delta(\mathbf{k}_1) = \mathcal{Z}^{(1)}(\mathbf{k}_1) \mathcal{Z}^{(1)}(-\mathbf{k}_1) P(k_1)$$

$$B_\Delta(\mathbf{k}_1, \mathbf{k}_2) = 2 \mathcal{Z}^{(1)}(\mathbf{k}_1) \mathcal{Z}^{(1)}(\mathbf{k}_2) \mathcal{Z}^{(2)}(\mathbf{k}_1, \mathbf{k}_2) P(k_1) P(k_2) + 2 \circlearrowleft_{\mathbf{k}_i}$$

# Galaxy clustering



- Galaxy clustering **kernels** in redshift space

$$\mathcal{Z}^{(1)}(\mathbf{k}_1) = b_1 + f \mu_1^2$$

$$\mathcal{Z}^{(2)}(\mathbf{k}_1, \mathbf{k}_2) = \frac{b_2}{2} + b_1 F_2(\mathbf{k}_1, \mathbf{k}_2) + b_{G_2} S_2(\mathbf{k}_1, \mathbf{k}_2) + f \mu_{12}^2 G_2(\mathbf{k}_1, \mathbf{k}_2) + \frac{f}{2} \mu_{12} k_{12} \left[ \frac{\mu_1}{k_1} \mathcal{Z}^{(1)}(\mathbf{k}_2) + \frac{\mu_2}{k_2} \mathcal{Z}^{(1)}(\mathbf{k}_1) \right]$$

- ...from which

$$P_{\Delta}(\mathbf{k}_1) = \mathcal{Z}^{(1)}(\mathbf{k}_1) \mathcal{Z}^{(1)}(-\mathbf{k}_1) P(k_1)$$

$$B_{\Delta}(\mathbf{k}_1, \mathbf{k}_2) = 2 \mathcal{Z}^{(1)}(\mathbf{k}_1) \mathcal{Z}^{(1)}(\mathbf{k}_2) \mathcal{Z}^{(2)}(\mathbf{k}_1, \mathbf{k}_2) P(k_1) P(k_2) + 2 \odot_{\mathbf{k}_i}$$

# Galaxy clustering



- Galaxy clustering **kernels** in redshift space

$$\mathcal{Z}^{(1)}(\mathbf{k}_1) = b_1 + f \mu_1^2$$

$$\mathcal{Z}^{(2)}(\mathbf{k}_1, \mathbf{k}_2) = \frac{b_2}{2} + b_1 \underline{F_2(\mathbf{k}_1, \mathbf{k}_2)} + b_{G_2} \underline{S_2(\mathbf{k}_1, \mathbf{k}_2)} + f \mu_{12}^2 \underline{G_2(\mathbf{k}_1, \mathbf{k}_2)} + \frac{f}{2} \mu_{12} k_{12} \left[ \frac{\mu_1}{k_1} \mathcal{Z}^{(1)}(\mathbf{k}_2) + \frac{\mu_2}{k_2} \mathcal{Z}^{(1)}(\mathbf{k}_1) \right]$$

- ...from which

$$P_{\Delta}(\mathbf{k}_1) = \mathcal{Z}^{(1)}(\mathbf{k}_1) \mathcal{Z}^{(1)}(-\mathbf{k}_1) P(k_1)$$

$$B_{\Delta}(\mathbf{k}_1, \mathbf{k}_2) = 2 \mathcal{Z}^{(1)}(\mathbf{k}_1) \mathcal{Z}^{(1)}(\mathbf{k}_2) \mathcal{Z}^{(2)}(\mathbf{k}_1, \mathbf{k}_2) P(k_1) P(k_2) + 2 \odot_{\mathbf{k}_i}$$

# Galaxy clustering



- Galaxy clustering **kernels** in redshift space

$$\mathcal{Z}^{(1)}(\mathbf{k}_1) = b_1 + f \mu_1^2$$

$$\mathcal{Z}^{(2)}(\mathbf{k}_1, \mathbf{k}_2) = \frac{b_2}{2} + b_1 \mathcal{F}_2(\mathbf{k}_1, \mathbf{k}_2) + b_{G_2} S_2(\mathbf{k}_1, \mathbf{k}_2) + f \mu_{12}^2 G_2(\mathbf{k}_1, \mathbf{k}_2) + \frac{f}{2} \mu_{12} k_{12} \left[ \frac{\mu_1}{k_1} \mathcal{Z}^{(1)}(\mathbf{k}_2) + \frac{\mu_2}{k_2} \mathcal{Z}^{(1)}(\mathbf{k}_1) \right]$$

- ...from which

$$P_{\Delta}(\mathbf{k}_1) = \mathcal{Z}^{(1)}(\mathbf{k}_1) \mathcal{Z}^{(1)}(-\mathbf{k}_1) P(k_1)$$

$$B_{\Delta}(\mathbf{k}_1, \mathbf{k}_2) = 2 \mathcal{Z}^{(1)}(\mathbf{k}_1) \mathcal{Z}^{(1)}(\mathbf{k}_2) \mathcal{Z}^{(2)}(\mathbf{k}_1, \mathbf{k}_2) P(k_1) P(k_2) + 2 \odot_{\mathbf{k}_i}$$

# Peculiar velocities



UNIVERSITÀ  
DI TORINO

- Galaxies' (radial) peculiar velocities are inferred from various types of observations (SNeIa luminosity-distance fluctuations, Tully-Fisher or Faber-Jackson relations, ...)



# Peculiar velocities



UNIVERSITÀ  
DI TORINO

- Galaxies' (radial) peculiar velocities are inferred from various types of observations (SNeIa luminosity-distance fluctuations, Tully-Fisher or Faber-Jackson relations, ...)
- RSD still arise, because in most cases the underlying peculiar velocity field of matter (as traced by galaxies) is reconstructed by positioning galaxies in a 3D comoving grid constructed from their redshift-space position—same as in galaxy clustering

# Peculiar velocities



- Galaxies' (radial) peculiar velocities are inferred from various types of observations (SN Ia luminosity-distance fluctuations, Tully-Fisher or Faber-Jackson relations, ...)
- RSD still arise, because in most cases the underlying peculiar velocity field of matter (as traced by galaxies) is reconstructed by positioning galaxies in a 3D comoving grid constructed from their redshift-space position—same as in galaxy clustering
- Analogously to what done before, I computed peculiar-velocity kernels...

$$u^{(m)}(\mathbf{k}) = \int_{\mathbf{q}_1} \dots \int_{\mathbf{q}_m} \delta^{(1)}(\mathbf{q}_1) \dots \delta^{(1)}(\mathbf{q}_m) \mathcal{U}^{(m)}(\mathbf{q}_1, \dots, \mathbf{q}_m) (2\pi)^3 \delta_D(\mathbf{q}_{1\dots m} - \mathbf{k})$$

# Peculiar velocities



- ...from which

$$P_u(\mathbf{k}_1) = \mathcal{U}^{(1)}(\mathbf{k}_1) \mathcal{U}^{(1)}(-\mathbf{k}_1) P(k_1)$$

$$B_u(\mathbf{k}_1, \mathbf{k}_2) = 2\mathcal{U}^{(1)}(\mathbf{k}_1) \mathcal{U}^{(1)}(\mathbf{k}_2) \mathcal{U}^{(2)}(\mathbf{k}_1, \mathbf{k}_2) P(k_1) P(k_2) + 2\mathcal{O}(k_i)$$

# Peculiar velocities



- ...from which

$$P_u(\mathbf{k}_1) = \mathcal{U}^{(1)}(\mathbf{k}_1) \mathcal{U}^{(1)}(-\mathbf{k}_1) P(k_1)$$

$$B_u(\mathbf{k}_1, \mathbf{k}_2) = 2\mathcal{U}^{(1)}(\mathbf{k}_1) \mathcal{U}^{(1)}(\mathbf{k}_2) \mathcal{U}^{(2)}(\mathbf{k}_1, \mathbf{k}_2) P(k_1) P(k_2) + 2\mathcal{O}_{\mathbf{k}_i}$$

- The kernel for the **power spectrum** coincides with the literature

$$\mathcal{U}^{(1)}(\mathbf{k}_1) = -i \mathcal{H} f D \frac{\mu_1}{k_1}$$

# Peculiar velocities



- ...from which

$$P_u(\mathbf{k}_1) = \mathcal{U}^{(1)}(\mathbf{k}_1) \mathcal{U}^{(1)}(-\mathbf{k}_1) P(k_1)$$

$$B_u(\mathbf{k}_1, \mathbf{k}_2) = 2\mathcal{U}^{(1)}(\mathbf{k}_1) \mathcal{U}^{(1)}(\mathbf{k}_2) \mathcal{U}^{(2)}(\mathbf{k}_1, \mathbf{k}_2) P(k_1) P(k_2) + 2\mathcal{O}_{\mathbf{k}_i}$$

- The kernel for the **power spectrum** coincides with the literature

$$\mathcal{U}^{(1)}(\mathbf{k}_1) = -i \mathcal{H} f D \frac{\mu_1}{k_1}$$

- The kernel for the **bispectrum** is a novel result

$$\mathcal{U}^{(2)}(\mathbf{k}_1, \mathbf{k}_2) = \mathcal{U}^{(1)}(\mathbf{k}_{12}) D \left[ G_2(\mathbf{k}_1, \mathbf{k}_2) - \frac{3}{2} f \mu_1 \mu_2 \frac{(k_{12})^2}{k_1 k_2} \right]$$

# Peculiar velocities



- ...from which

$$P_u(\mathbf{k}_1) = \mathcal{U}^{(1)}(\mathbf{k}_1) \mathcal{U}^{(1)}(-\mathbf{k}_1) P(k_1)$$

$$B_u(\mathbf{k}_1, \mathbf{k}_2) = 2\mathcal{U}^{(1)}(\mathbf{k}_1) \mathcal{U}^{(1)}(\mathbf{k}_2) \mathcal{U}^{(2)}(\mathbf{k}_1, \mathbf{k}_2) P(k_1) P(k_2) + 2\mathcal{O}_{\mathbf{k}_i}$$

- The kernel for the **power spectrum** coincides with the literature

$$\mathcal{U}^{(1)}(\mathbf{k}_1) = -i \mathcal{H} f D \frac{\mu_1}{k_1}$$

- The kernel for the **bispectrum** is a novel result

$$\mathcal{U}^{(2)}(\mathbf{k}_1, \mathbf{k}_2) = \mathcal{U}^{(1)}(\mathbf{k}_{12}) D \left[ \underline{G_2(\mathbf{k}_1, \mathbf{k}_2)} - \frac{3}{2} f \mu_1 \mu_2 \frac{(k_{12})^2}{k_1 k_2} \right]$$

# Momentum density



- Momentum density is the **density-weighted** peculiar velocity field

$$p := (1 + \Delta) u$$

- ...from which

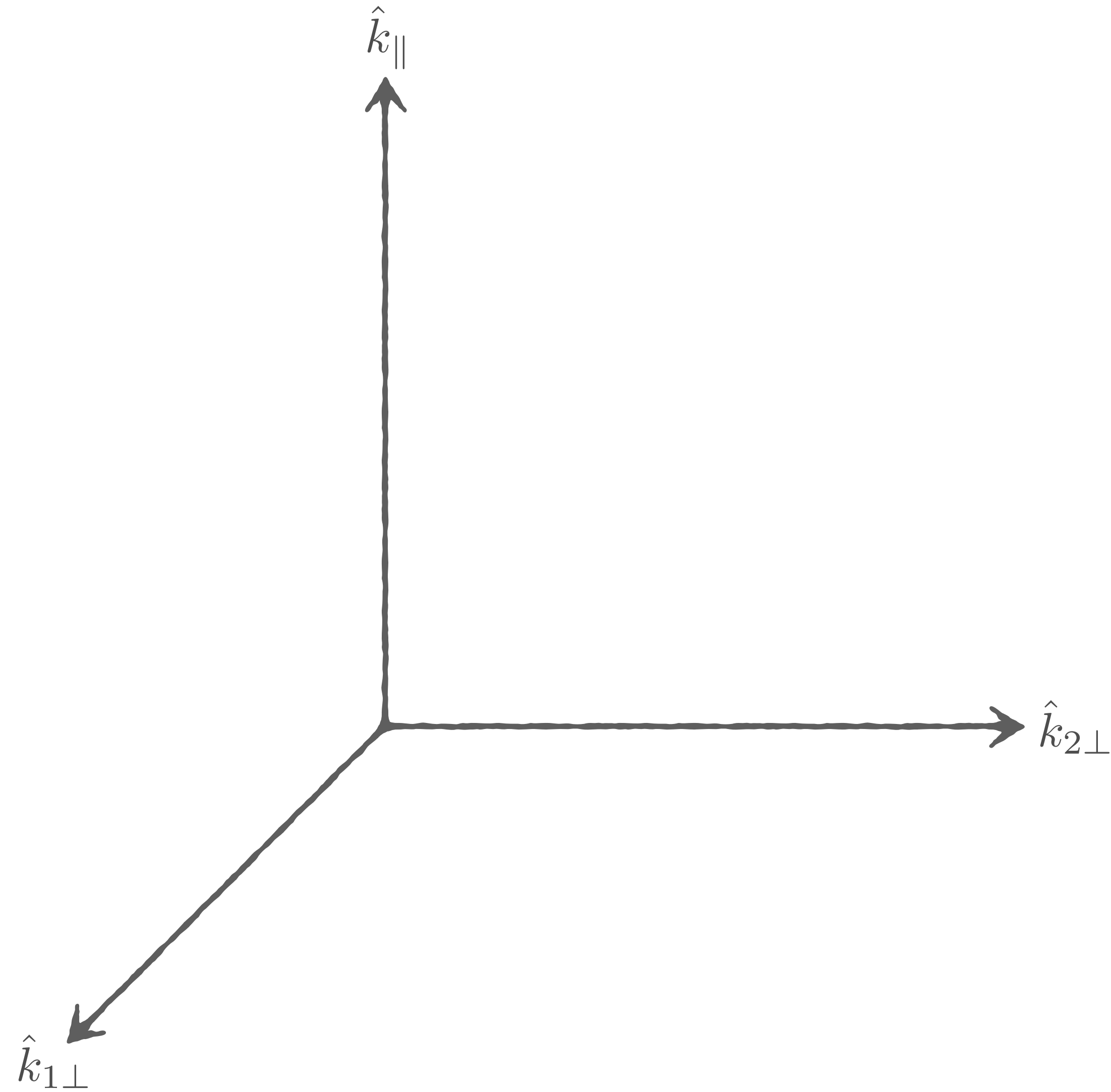
$$\mathcal{P}^{(1)}(\mathbf{k}_1) = \mathcal{U}^{(1)}(\mathbf{k}_1)$$

$$\mathcal{P}^{(2)}(\mathbf{k}_1, \mathbf{k}_2) = \mathcal{U}^{(2)}(\mathbf{k}_1, \mathbf{k}_2) + \mathcal{U}^{(1)}(\mathbf{k}_{12}) D \left[ \frac{f}{2} \mu_1 \mu_2 \frac{(k_{12})^2}{k_1 k_2} + \frac{b_1}{2} \left( \frac{\mu_1}{k_1} + \frac{\mu_2}{k_2} \right) \frac{k_{12}}{\mu_{12}} \right]$$

# Constructing bispectra



UNIVERSITÀ  
DI TORINO

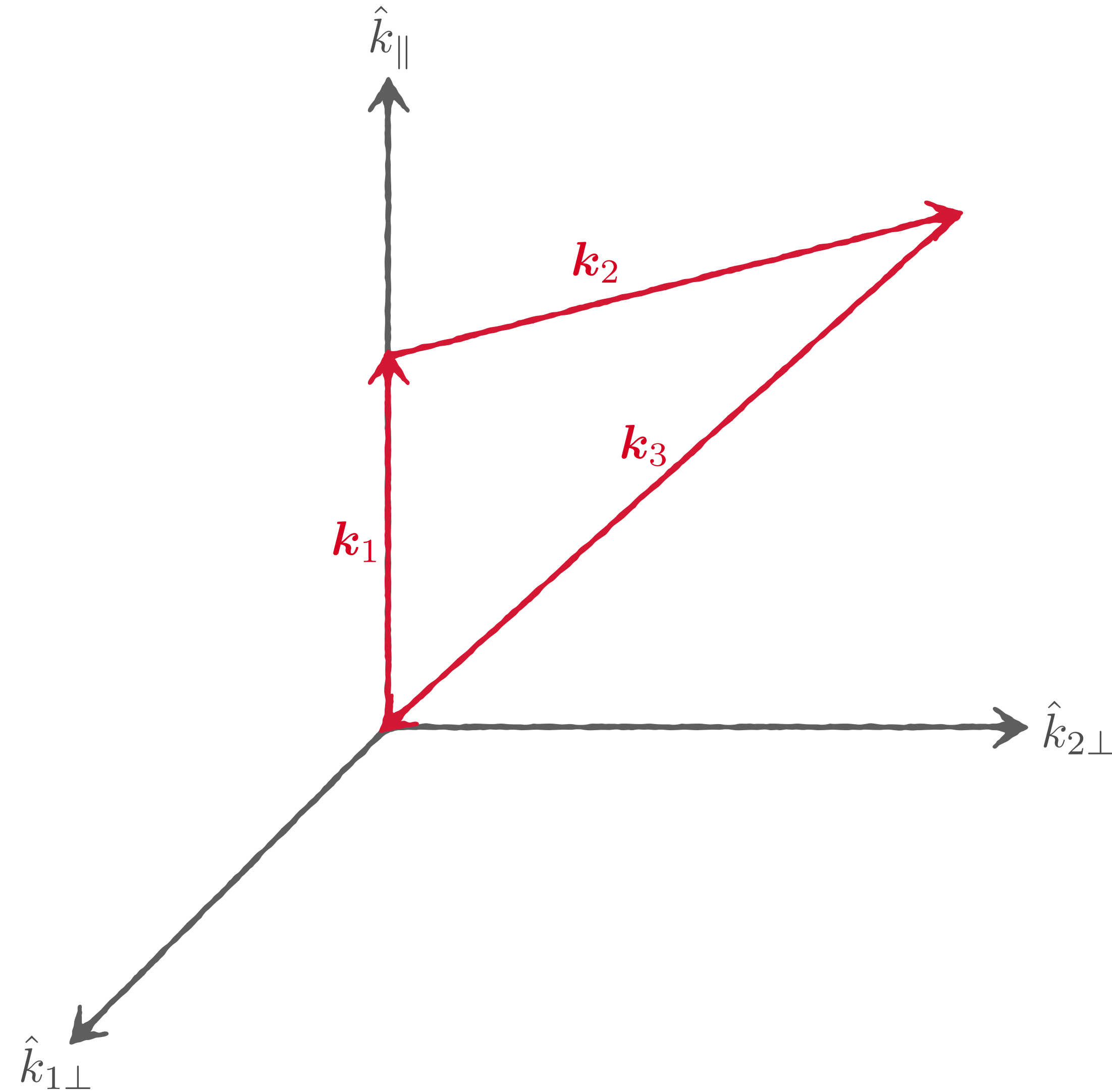




# Constructing bispectra



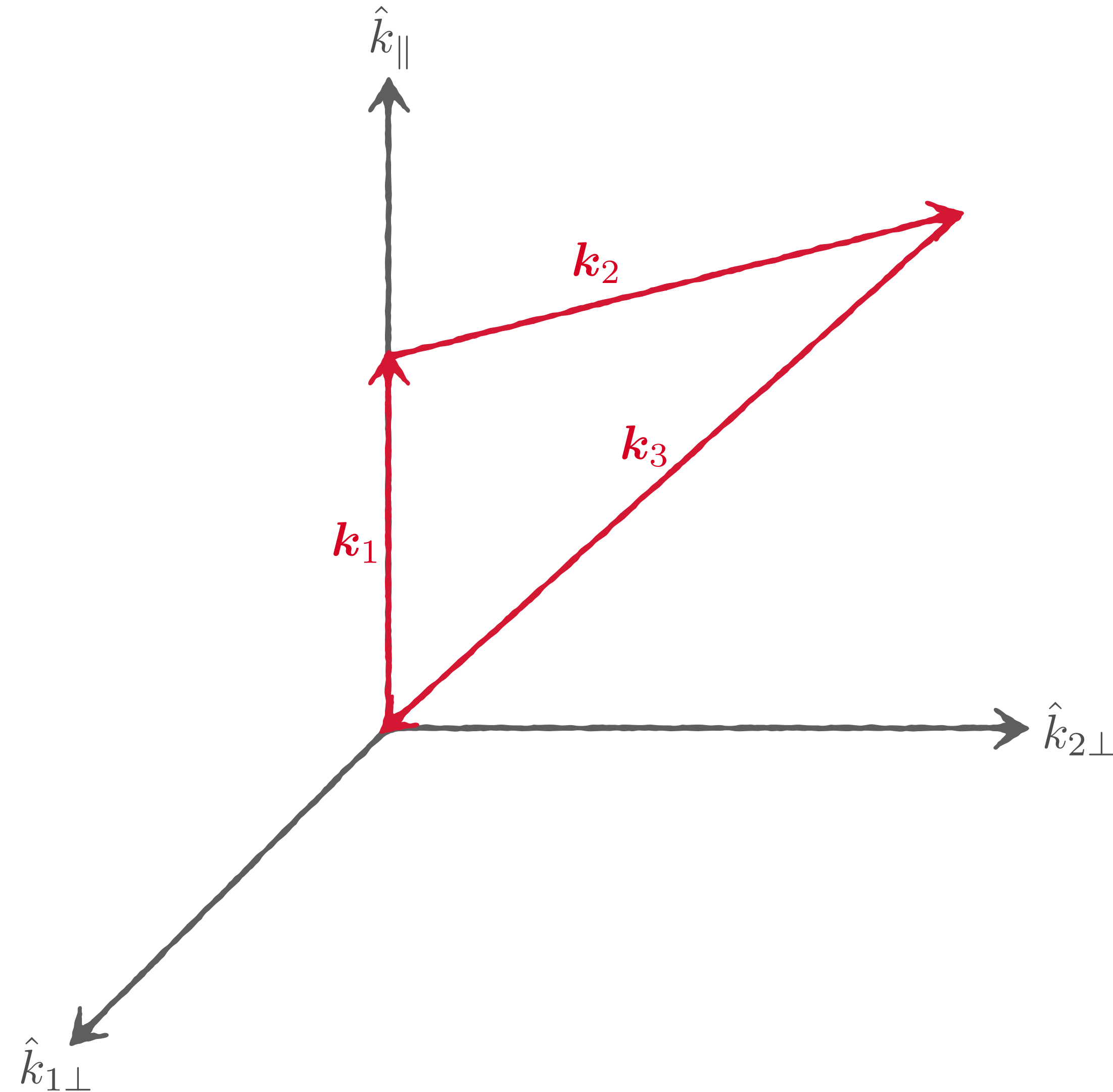
UNIVERSITÀ  
DI TORINO



# Constructing bispectra



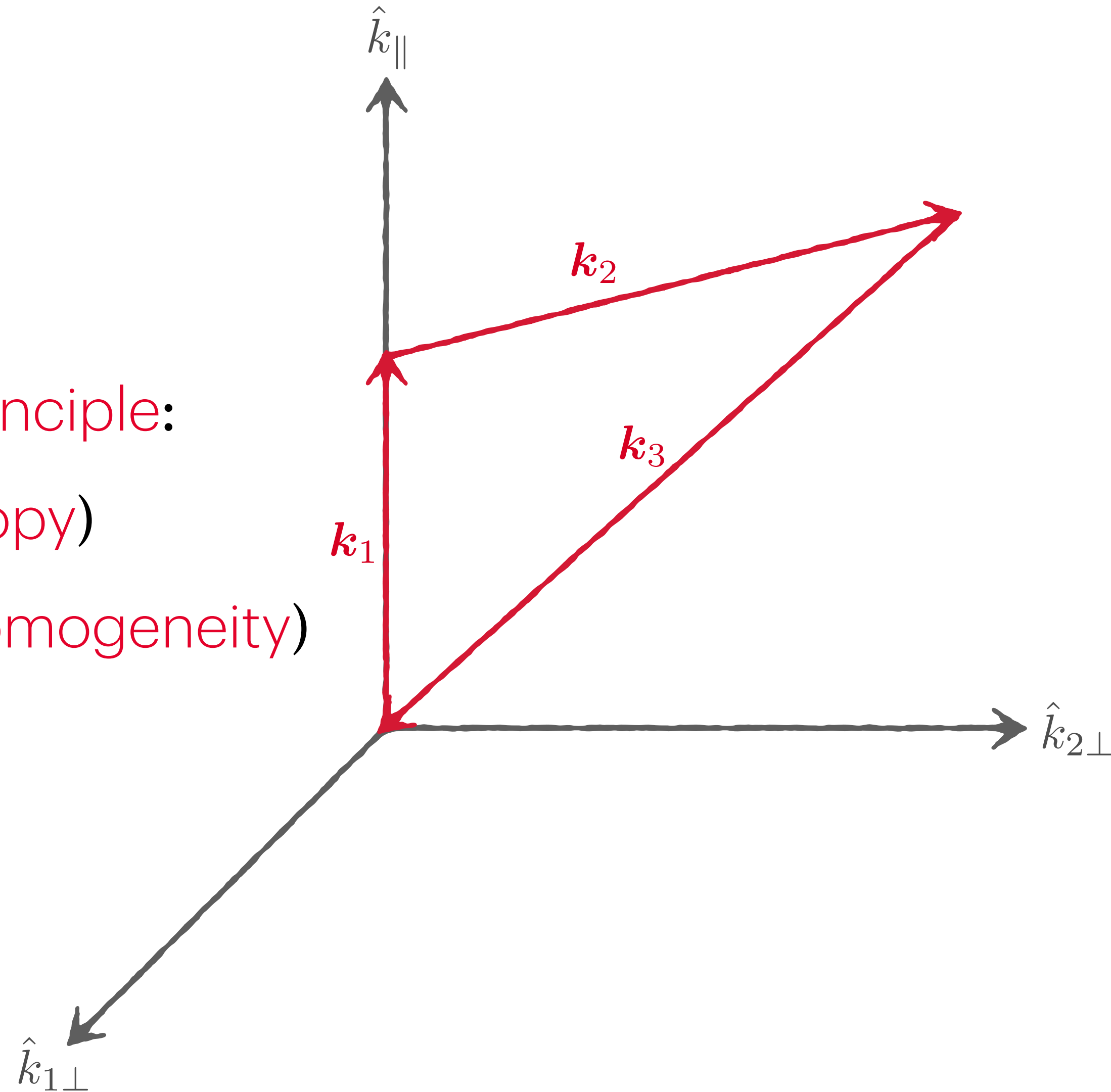
- # of dof:
  - 3 3D  $k_i = 9$



# Constructing bispectra



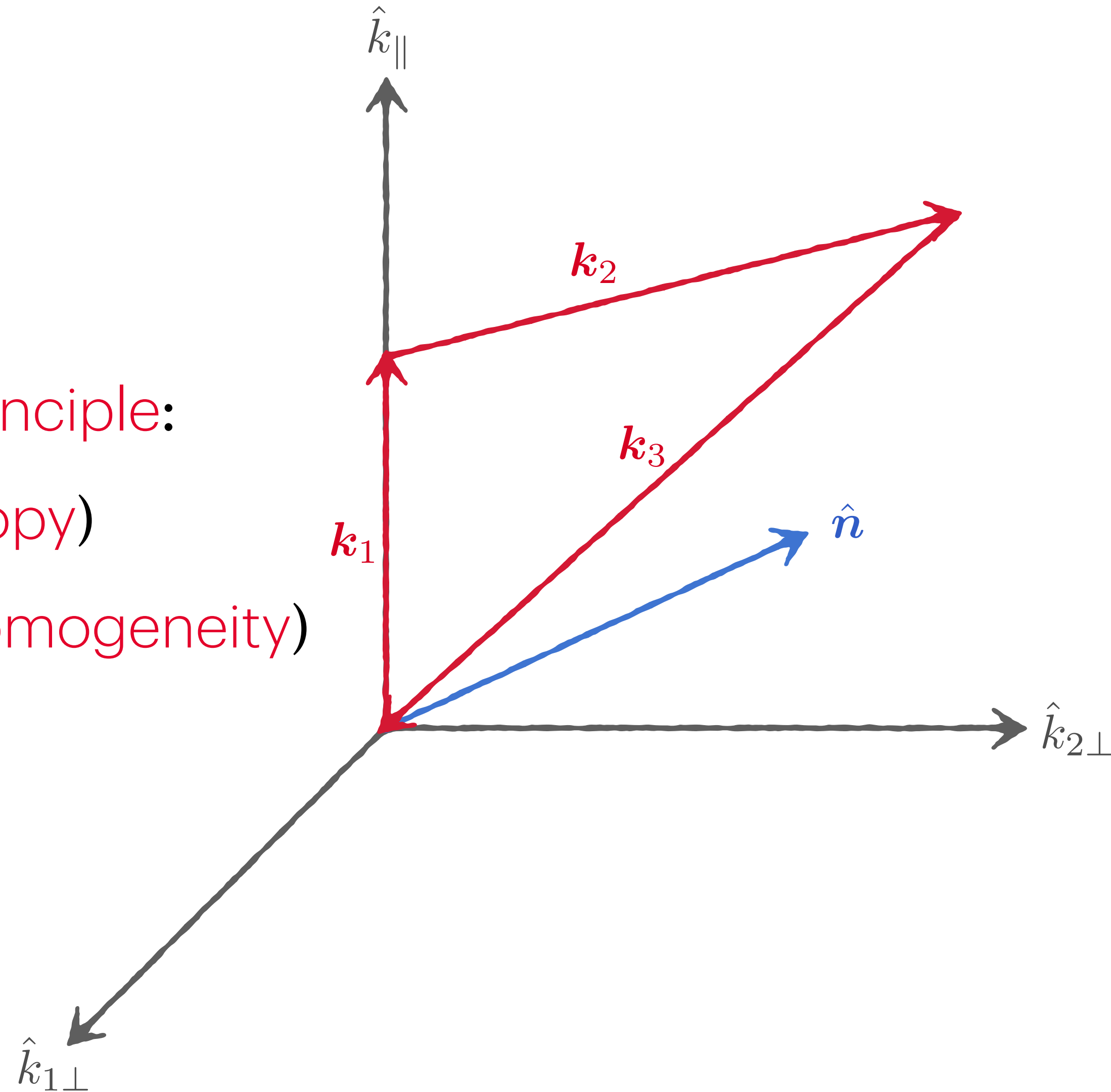
- # of dof:
  - 3 3D  $k_i = 9$
- But cosmological principle:
  - -3 rotations (isotropy)
  - -3 translations (homogeneity)



# Constructing bispectra



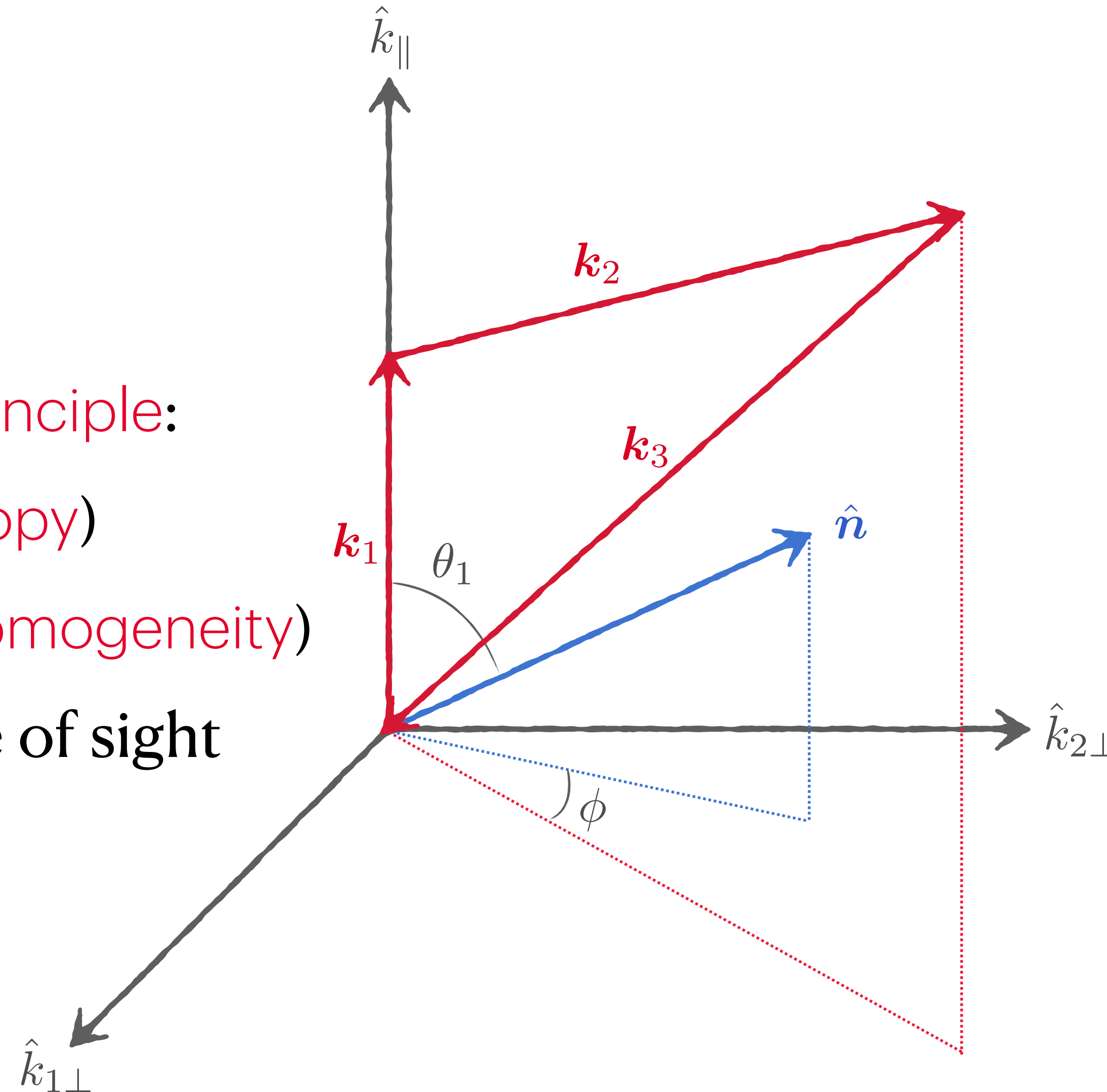
- # of dof:
  - 3 3D  $k_i = 9$
- But cosmological principle:
  - -3 rotations (isotropy)
  - -3 translations (homogeneity)



# Constructing bispectra



- # of dof:
  - 3 3D  $k_i = 9$
- But cosmological principle:
  - -3 rotations (isotropy)
  - -3 translations (homogeneity)
  - +2 angles w.r.t. line of sight
- 5 dof

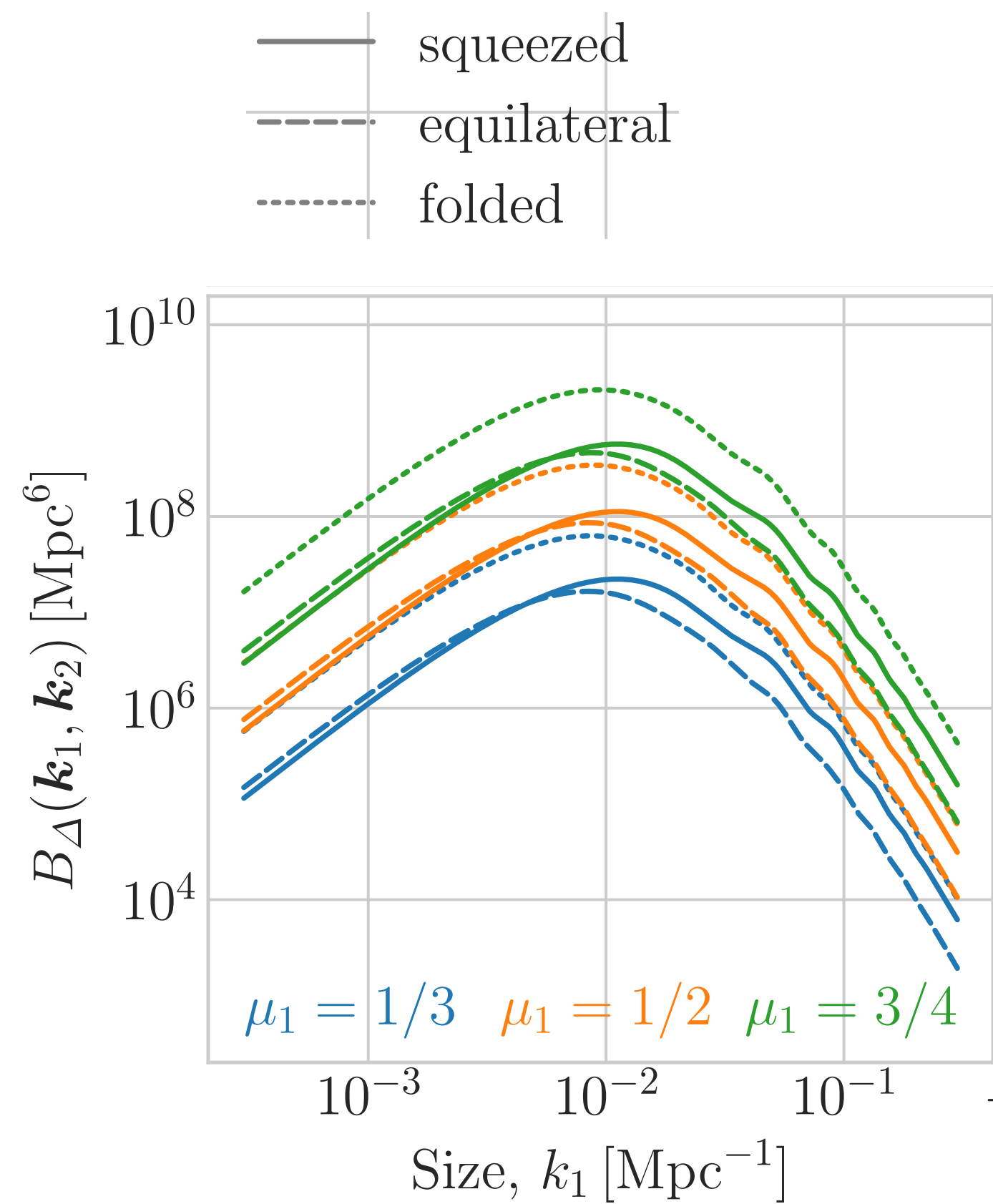


# Comparing bispectra

[SC 2024 (in prep.)]



UNIVERSITÀ  
DI TORINO

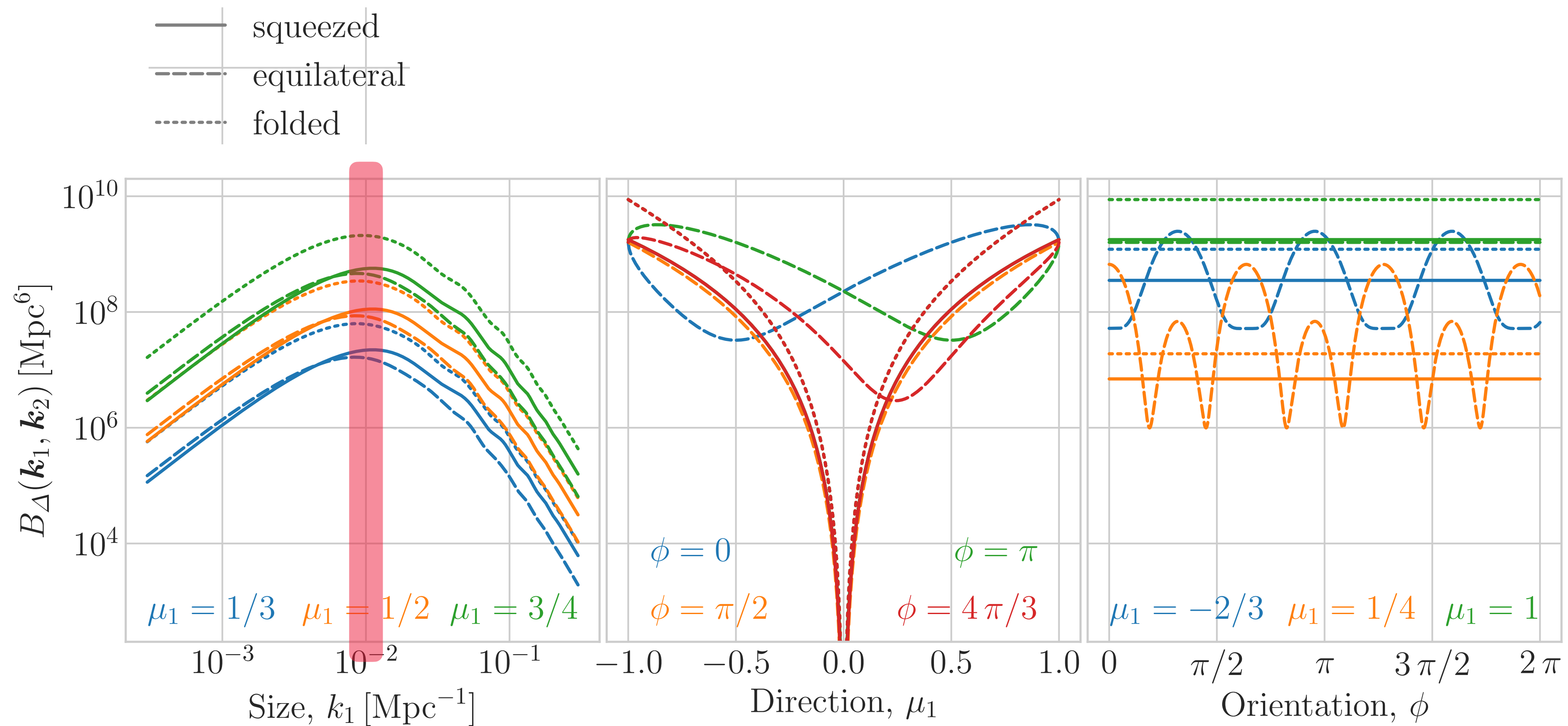


# Comparing bispectra



[SC 2024 (in prep.)]

UNIVERSITÀ  
DI TORINO

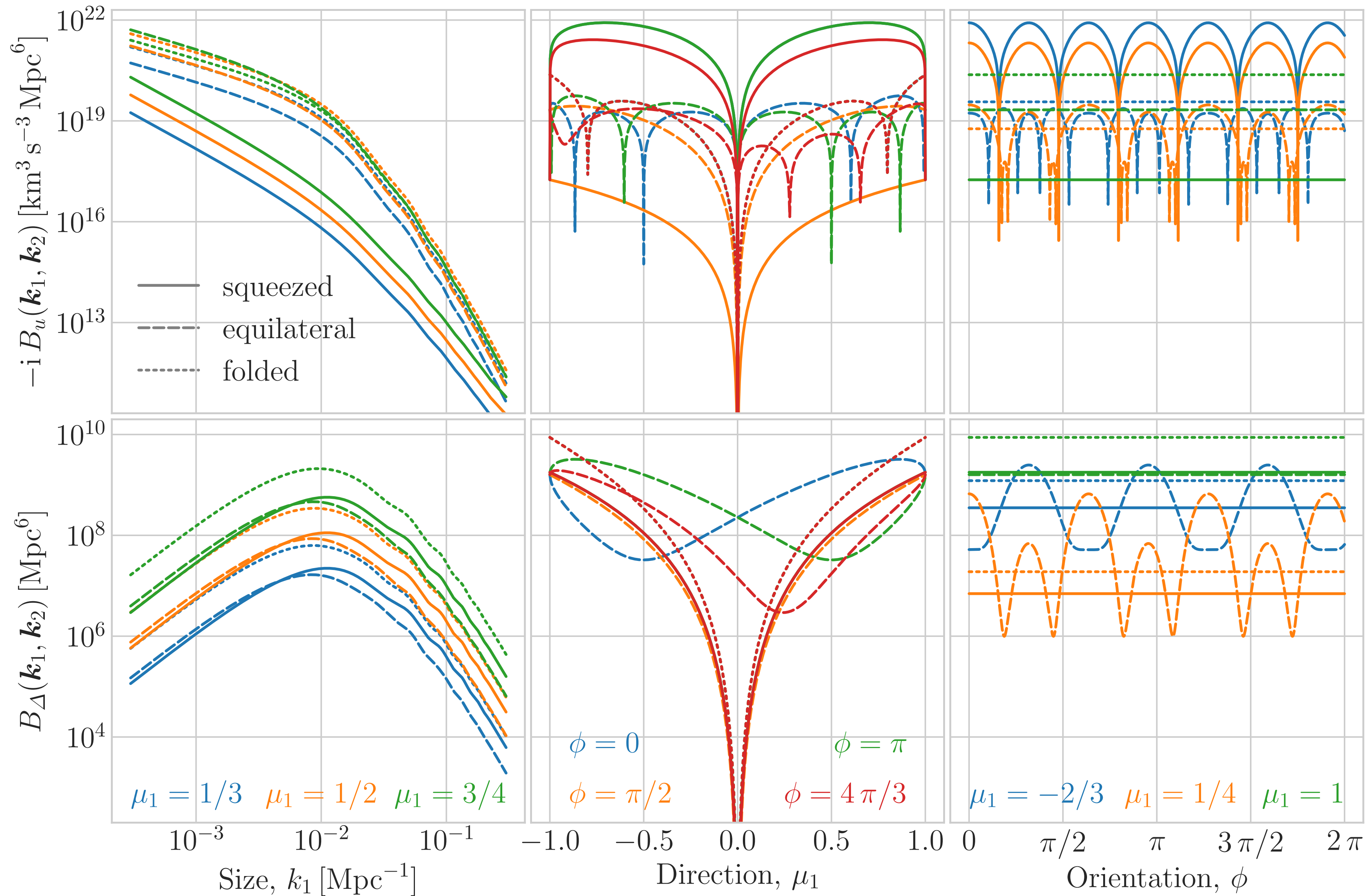


# Comparing bispectra

[SC 2024 (in prep.)]



UNIVERSITÀ  
DI TORINO





# Detectability



$$\tilde{P}_X = P_X + \begin{cases} 1/\bar{n}_g(\bar{z}_i) & \text{if } X = \Delta \\ \sigma_v^2/\bar{n}_v(\bar{z}_i) & \text{if } X = \{u, p\} \end{cases}$$

# Detectability



$$\tilde{P}_X = \textcircled{P_X} + \begin{cases} 1/\bar{n}_g(\bar{z}_i) & \text{if } X = \Delta \\ \sigma_v^2/\bar{n}_v(\bar{z}_i) & \text{if } X = \{u, p\} \end{cases}$$

# Detectability



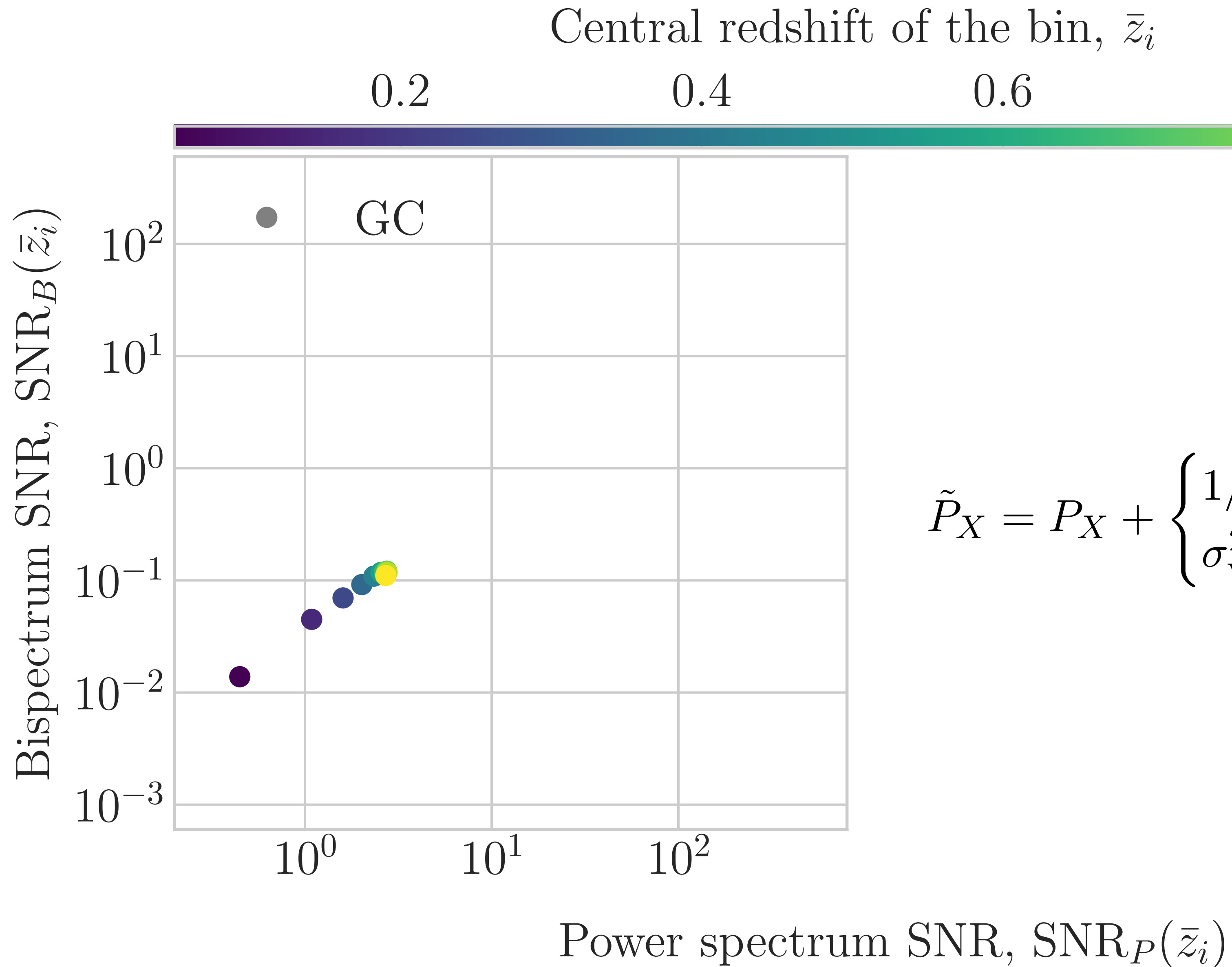
$$\tilde{P}_X = \underbrace{P_X}_{\text{red circle}} + \begin{cases} 1/\bar{n}_g(\bar{z}_i) & \text{if } X = \Delta \\ \sigma_v^2/\bar{n}_v(\bar{z}_i) & \text{if } X = \{u, p\} \end{cases}$$

# Detectability



UNIVERSITÀ  
DI TORINO

[SC 2024 (in prep.)]

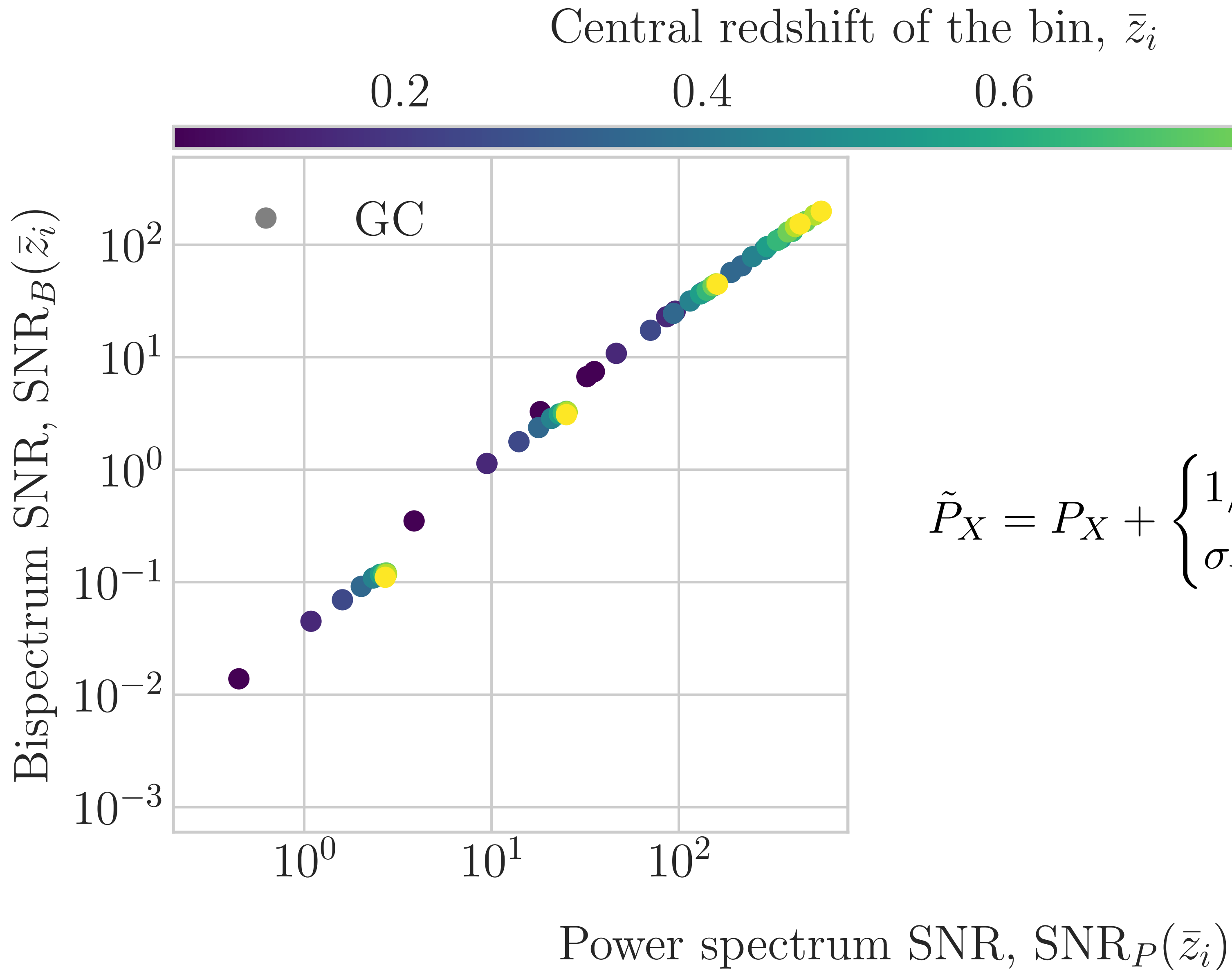


$$\tilde{P}_X = P_X + \begin{cases} 1/\bar{n}_g(\bar{z}_i) & \text{if } X = \Delta \\ \sigma_v^2/\bar{n}_v(\bar{z}_i) & \text{if } X = \{u, p\} \end{cases}$$

# Detectability



[SC 2024 (in prep.)]

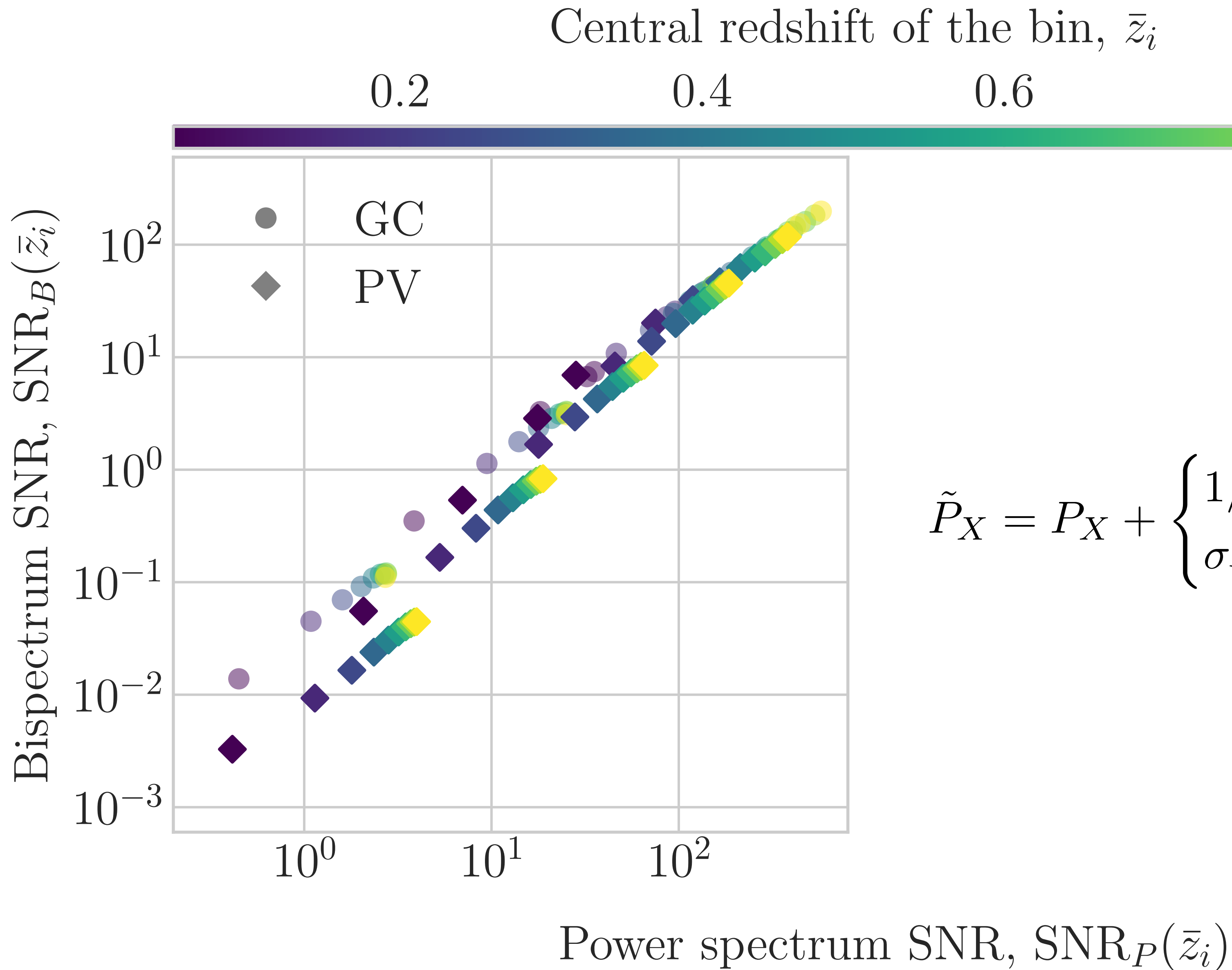


$$\tilde{P}_X = P_X + \begin{cases} 1/\bar{n}_g(\bar{z}_i) & \text{if } X = \Delta \\ \sigma_v^2/\bar{n}_v(\bar{z}_i) & \text{if } X = \{u, p\} \end{cases}$$

# Detectability



[SC 2024 (in prep.)]

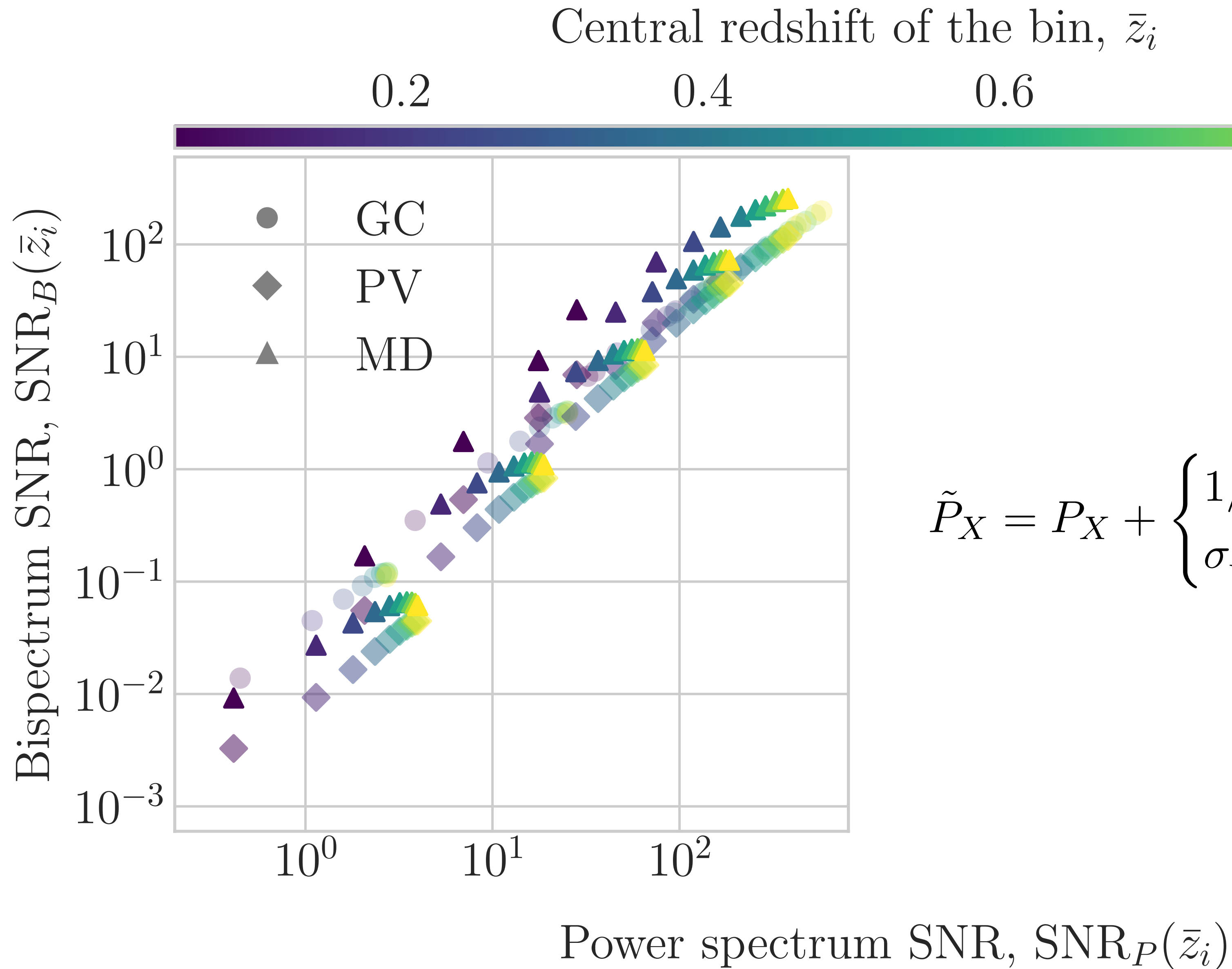


$$\tilde{P}_X = P_X + \begin{cases} 1/\bar{n}_g(\bar{z}_i) & \text{if } X = \Delta \\ \sigma_v^2/\bar{n}_v(\bar{z}_i) & \text{if } X = \{u, p\} \end{cases}$$

# Detectability



[SC 2024 (in prep.)]



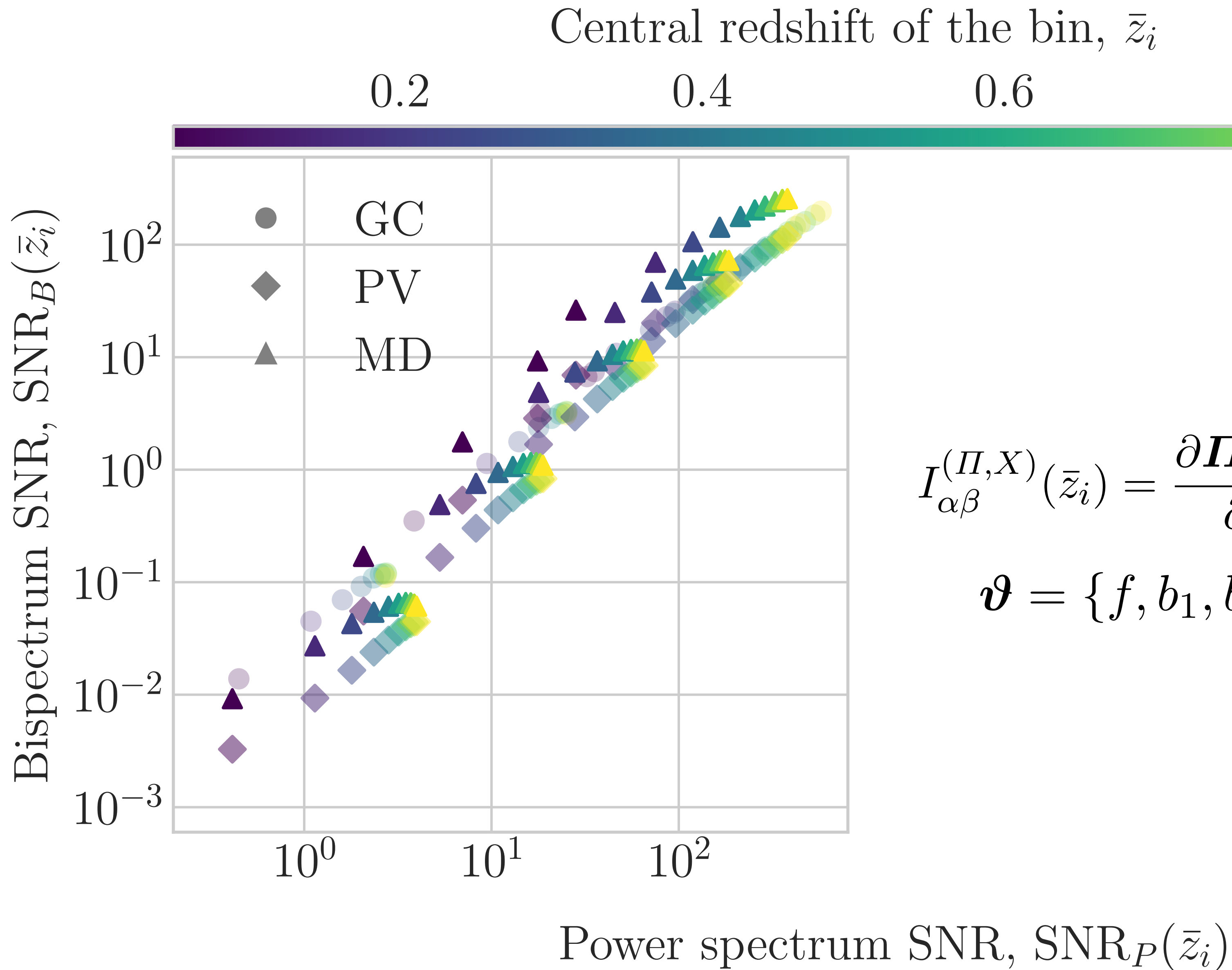
$$\tilde{P}_X = P_X + \begin{cases} 1/\bar{n}_g(\bar{z}_i) & \text{if } X = \Delta \\ \sigma_v^2/\bar{n}_v(\bar{z}_i) & \text{if } X = \{u, p\} \end{cases}$$

# Information content



UNIVERSITÀ  
DI TORINO

[SC 2024 (in prep.)]



$$I_{\alpha\beta}^{(\Pi, X)}(\bar{z}_i) = \frac{\partial \Pi_X^H(\bar{z}_i)}{\partial \vartheta_\alpha} \mathbf{C}^{-1}(\bar{z}_i) \frac{\partial \Pi_X(\bar{z}_i)}{\partial \vartheta_\beta}$$

$$\vartheta = \{f, b_1, b_2, b_{\mathcal{G}_2}, P_{\text{shot}}, B_{\text{shot}}\}$$

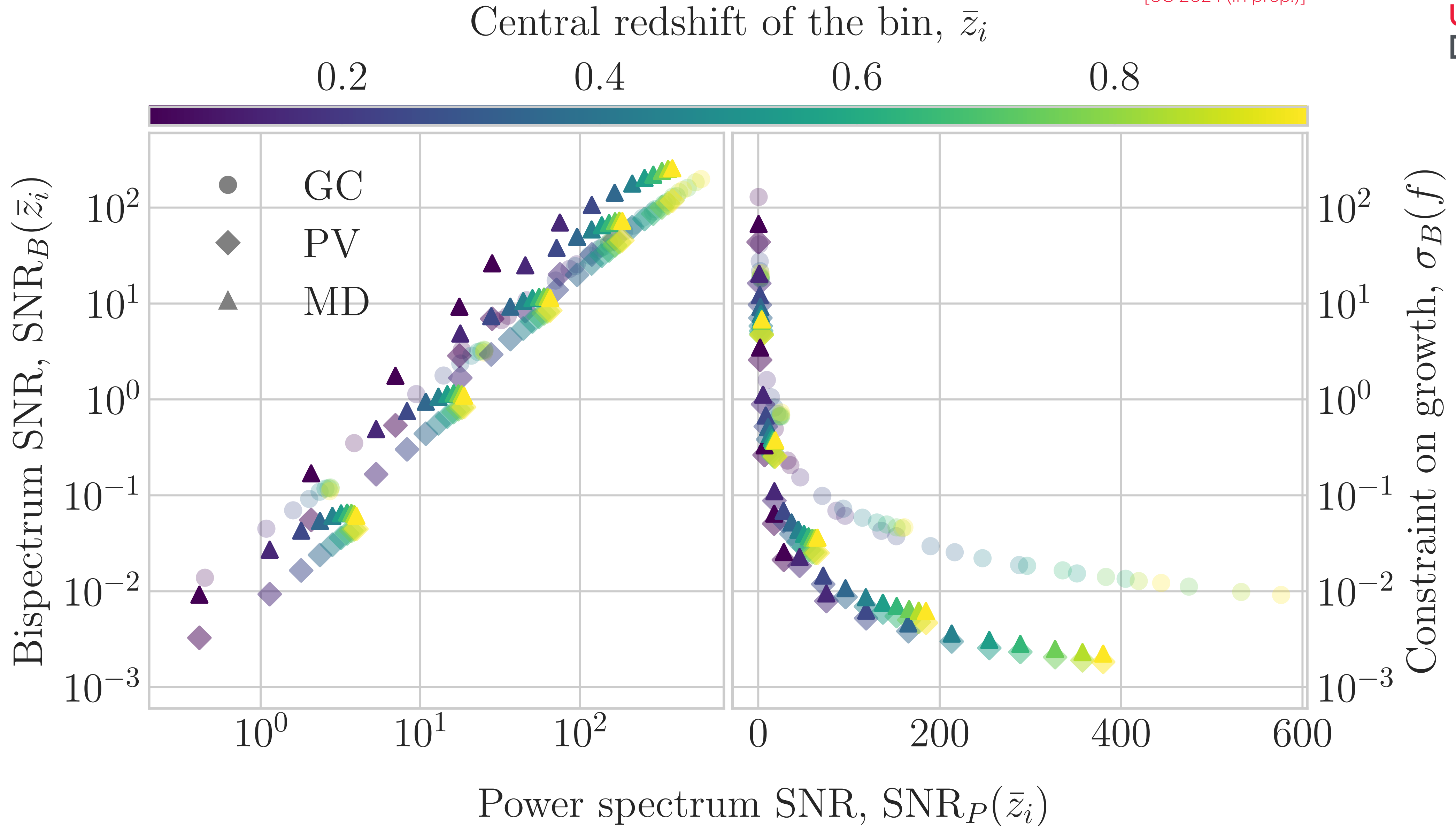


# Information content



UNIVERSITÀ  
DI TORINO

[SC 2024 (in prep.)]



# Information content



UNIVERSITÀ  
DI TORINO

[SC 2024 (in prep.)]

